

Bose-Hubbard models: Phases & Excitations

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**A tribute to Professor Satyendra Nath Bose on his 125th birth anniversary
International Workshop on Bose-Einstein Condensation and related
phenomena (IWBECP)
26-28th March, 2018**

Outline

Bose-Hubbard Models:

- Physical Realizations.
- Mean-field Theory : homogeneous
- Mean-field Theory : inhomogeneous
- Random Phase Approximation for Excitations.
- Finite Temperature:
- Bose Mixtures: Phases and Excitations
- Conclusions.

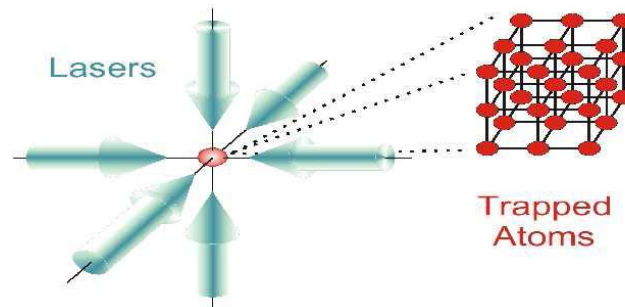
Physical Realizations

Interacting Bose Systems:

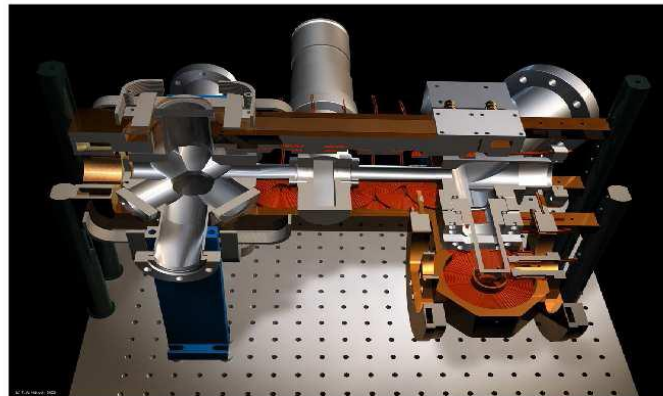
- ^4He in vycor or aerogel (disorder)
- Microfabricated Josephson junction arrays.
- Disorder-driven superconductor-insulator transition (e.g., thin films of bismuth).
- Type II superconductors with columnar defects.
- Ultra cold atoms in optical potentials.

Ultra Cold Atoms in Optical Lattices

Interference of standing wave laser beams is used to trap atoms in a periodic lattice potential:

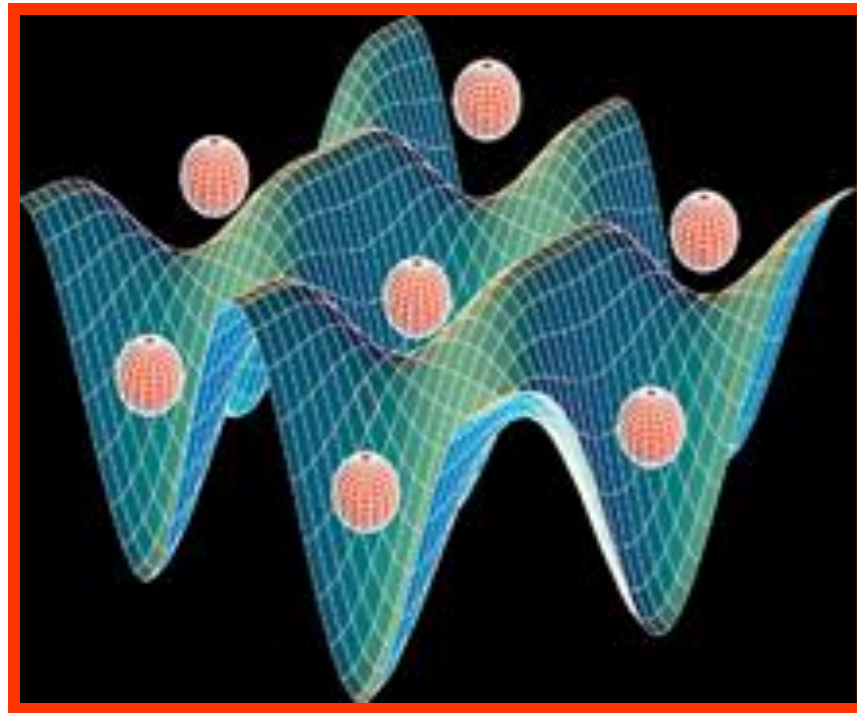


Munich: I. Bloch, T. Haensch et al.

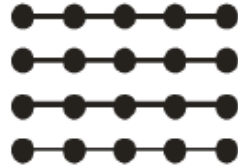


Cold Atoms in Optical Lattices

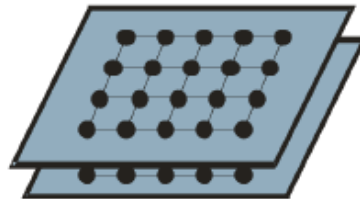
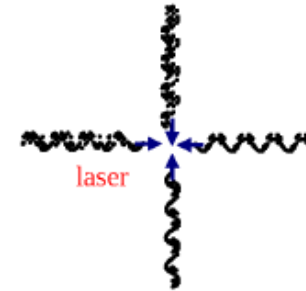
Interference of standing wave laser beams is used to trap atoms in a periodic lattice potential:



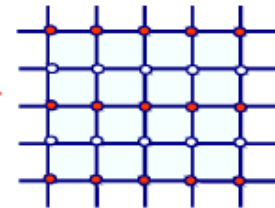
Lattice designs



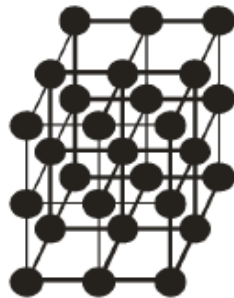
1D



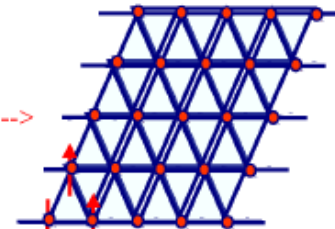
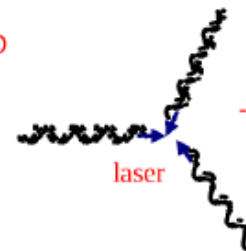
2D different internal states ----->



square lattice

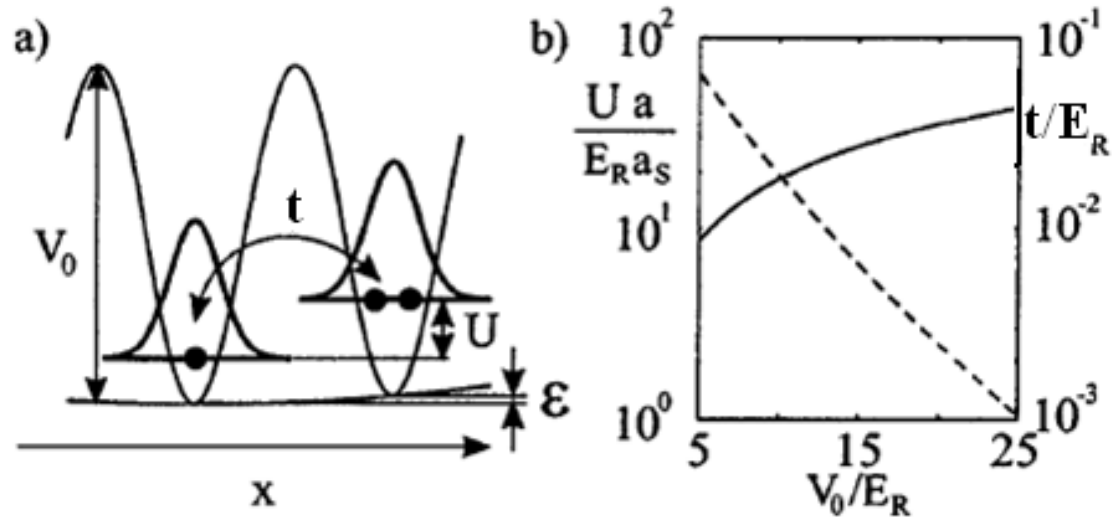


3D



triangular lattice

Model



$$\frac{t}{E_r} = \frac{4}{\sqrt{\pi}} \left(\frac{V_0}{E_r} \right)^{3/4} e^{-2\sqrt{\frac{V_0}{E_r}}}$$

$$E_r = \frac{\hbar^2 k^2}{2m} \quad \text{Recoil energy}$$

$$\frac{U}{E_r} = \sqrt{8\pi} \frac{a_s}{a} \left(\frac{V_0}{E_r} \right)^{3/4}$$

a_s s-wave scattering length

$a = \lambda/2$ Lattice constant

Jaksch et. al. (1998) PRL **81**, 3108

Bose-Hubbard model

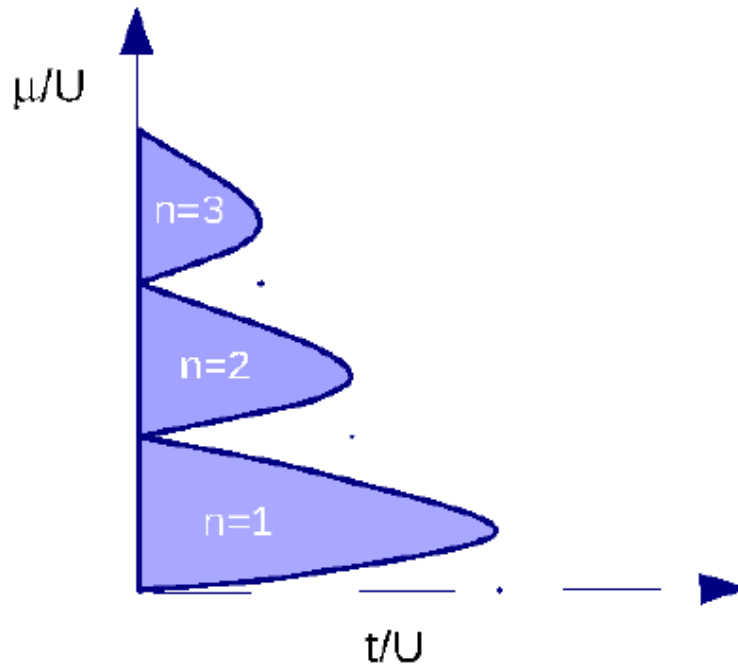
$$H = -t \sum_{\langle i,j \rangle} (a_i^+ a_j + h.c) + \frac{U}{2} \sum_i n_i (n_i - 1) - \mu \sum_i n_i + V_T \sum_i n_i (\vec{r}_i - \vec{r}_0)^2$$

Where V_T is the strength of the trap potential and \vec{r}_0 is the center of the lattice

$$|\Psi_{SF}\rangle \xrightarrow{\frac{U}{t} \approx 0} \propto \left(\sum_i a_i^+ \right)^N |0\rangle$$

$$|\Psi_{MI}\rangle \xrightarrow{t \approx 0} \propto \prod_i (a_i^+)^N |0\rangle$$

Bose-Hubbard model: Phase Diagram



- M. Fisher et. al. Phys. Rev. B. **40**, 546 (1989)
- Sheshadri et. al. Europhys. Lett. **22** 257 (1993)
- Experiments: M. Greiner et. al. Nature **415**, 39 (2002).

Bose-Hubbard model: Mean-field theory

$$H = -t_a \sum_{\langle i,j \rangle} (a_i^+ a_j + h.c) + \frac{U}{2} \sum_i n_i (n_i - 1) - \mu \sum_i n_i$$

$$a_i^+ a_j \approx \langle a_i^+ \rangle a_j + a_i^+ \langle a_j \rangle - \langle a_i^+ \rangle \langle a_j \rangle$$

Superfluid order parameter $\psi_i^a = \langle a_i \rangle$

$$H = \sum_i H_i^{MF} - t \sum_{\langle i,j \rangle} (\delta a_i^+ \delta a_j + h.c)$$

$$H_i^{MF} = \frac{U}{2} \hat{n}_i (\hat{n}_i - 1) - \mu n_i - (\psi_i^* a_i + \psi_i a_i^+) + |\psi_i|^2$$

Sheshadri et. al. Europhys. Lett. **22** 257 (1993)

Bose-Hubbard model

Mean-field theory

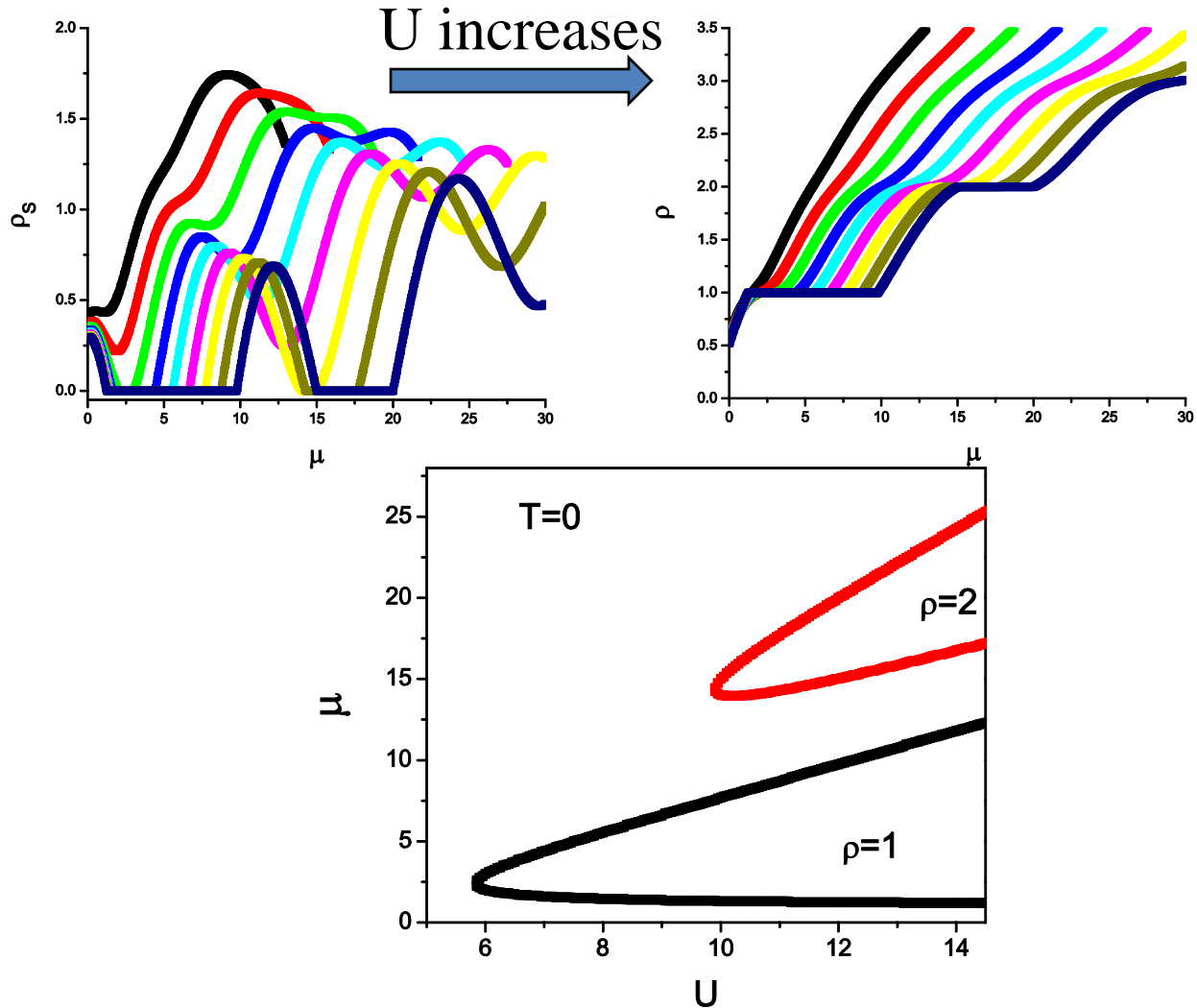
$$H_i^{MF} = \frac{U}{2} \hat{n}_i (\hat{n}_i - 1) - \mu n_i - (\psi_i^* a_i + \psi_i a_i^+) + |\psi_i|^2$$

- Superfluid density $(\rho_s)_i = |\psi_i|^2$
- Density $\rho_i = \langle \hat{n}_i \rangle$
- Compressibility $\kappa_i = \frac{\partial \rho_i}{\partial \mu_i}$

SF Phase: $\rho_s > 0$ $\kappa > 0$

MI Phase: $\rho_s = 0$ $\kappa = 0$

Bose-Hubbard model: Phase Diagram



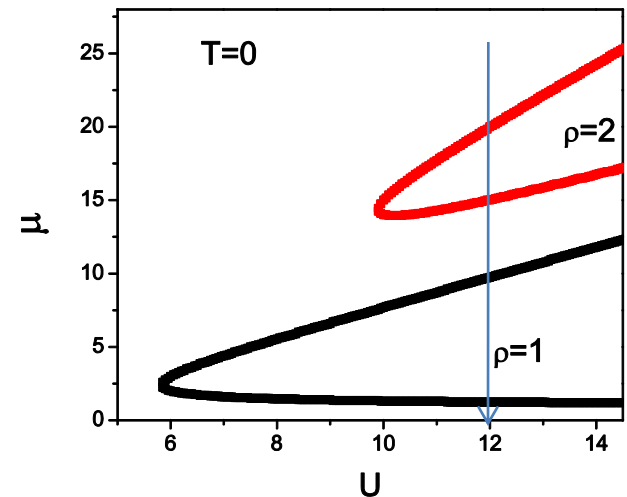
Bose-Hubbard model: Inhomogeneous MFT

$$H^{MF} = \sum_i H_i + \frac{1}{2z} \sum_{\langle ij \rangle} (\psi_i^* \psi_j + \psi_i \psi_j^*)$$

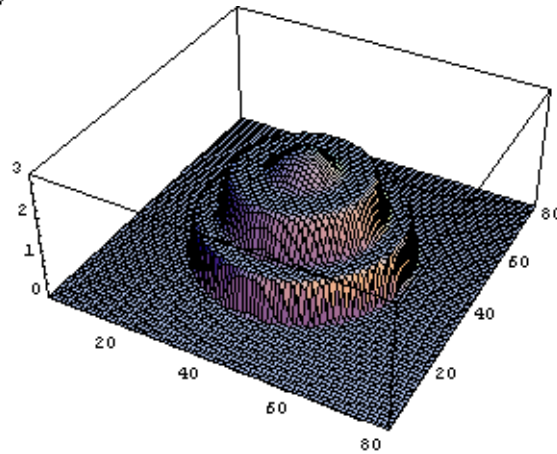
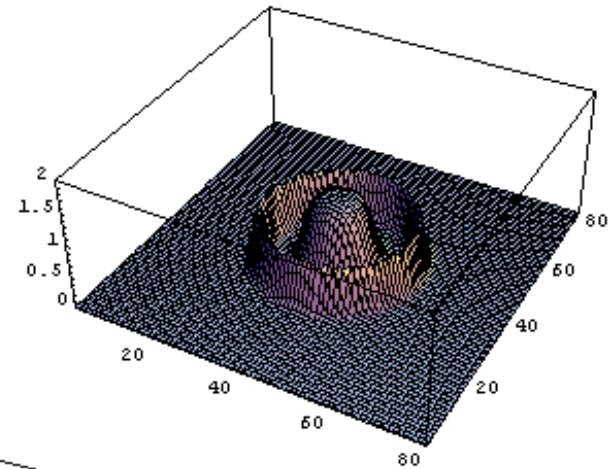
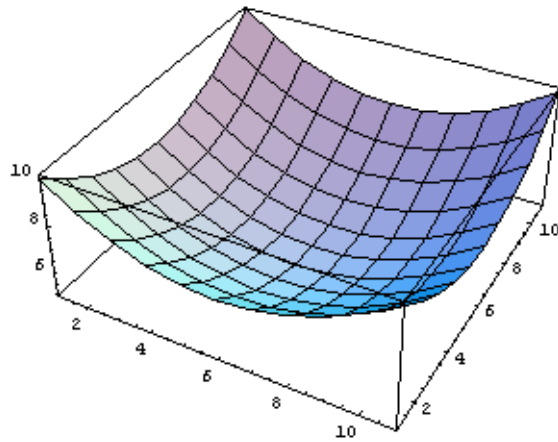
$$H_i = \frac{U}{2} \hat{n}_i (\hat{n}_i - 1) - \mu_i n_i - (\phi_i^* a_i + \phi_i a_i^+)$$

$$\phi_i = \frac{1}{z} \sum_{\delta} \psi_{i+\delta}$$

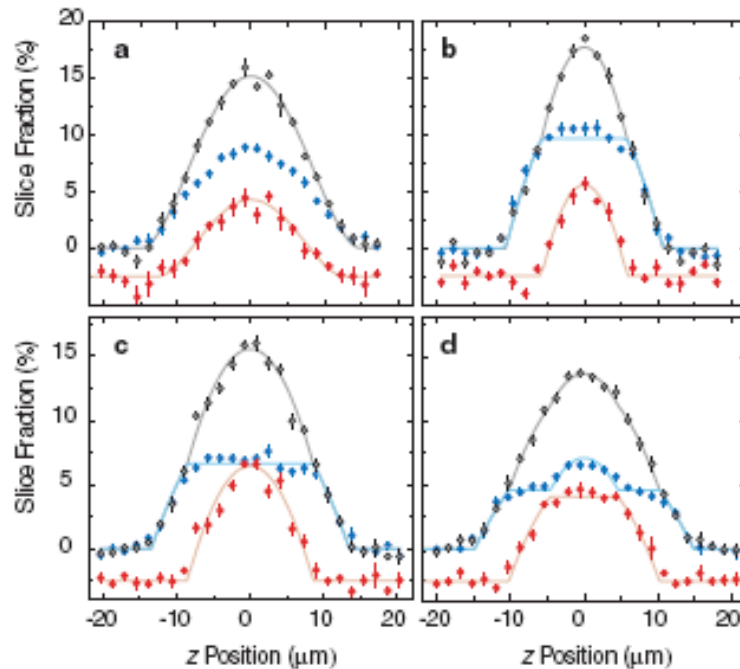
$$\mu_i = \mu - V_T (\vec{r}_i - \vec{r}_0)^2$$



Bose-Hubbard model: SF & MI shells



Bose-Hubbard model: SF & MI shells



	Number of atoms	V_0 (E_r)
a	1.0×10^5	3
b	1.0×10^5	22
c	2.0×10^5	22
d	3.5×10^5	22

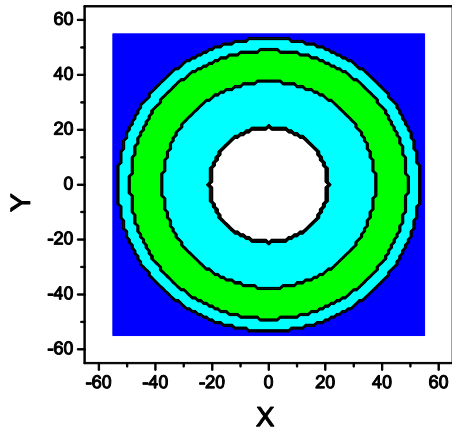
Gray: Total density distribution

Red: Density distribution with $\rho=2$

Blue: Gray - Red

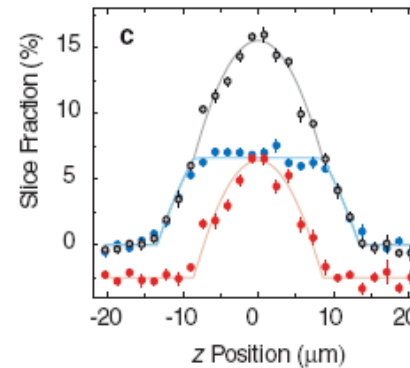
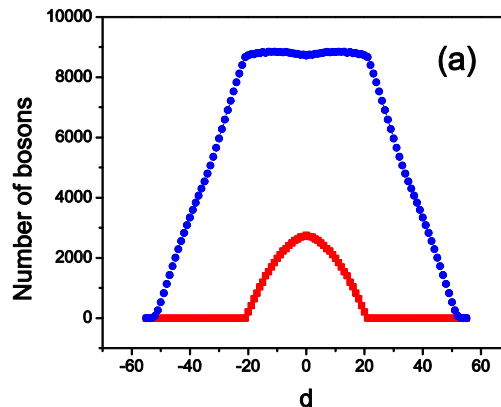
Folling et. al. PRL **97** 060403 (2006)

Bose-Hubbard model: Density Distribution & Shells: Result



SF
MI ($\rho=1$)
MI ($\rho=2$)

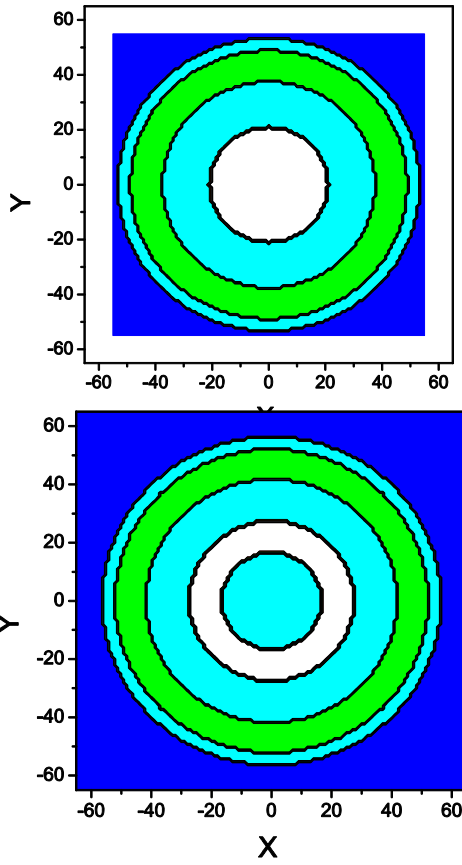
	Number of atoms	V_0 (E_r)
a	7.1×10^5	15



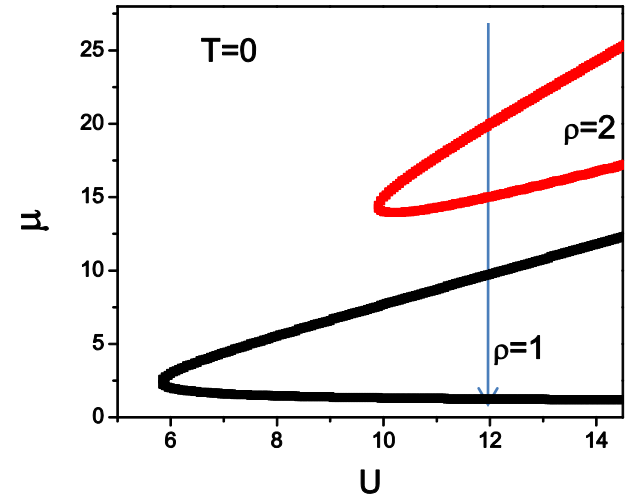
Red: Density distribution with $\rho=2$

Blue: Total Distribution - **Red** Phys. Rev. B **85**, 214524 (2012)

Bose-Hubbard model: Density Distribution & Shells: Result



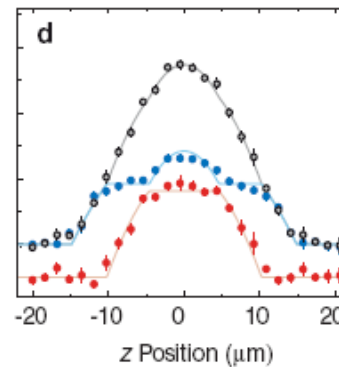
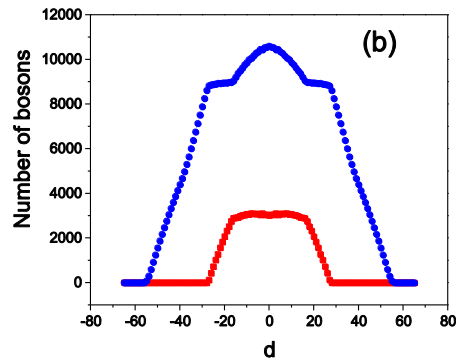
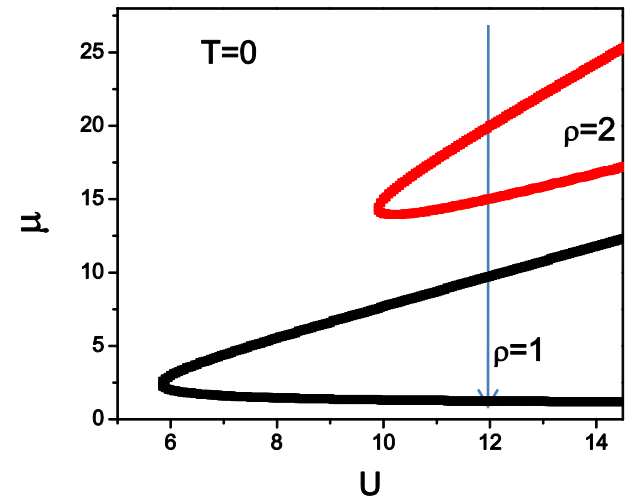
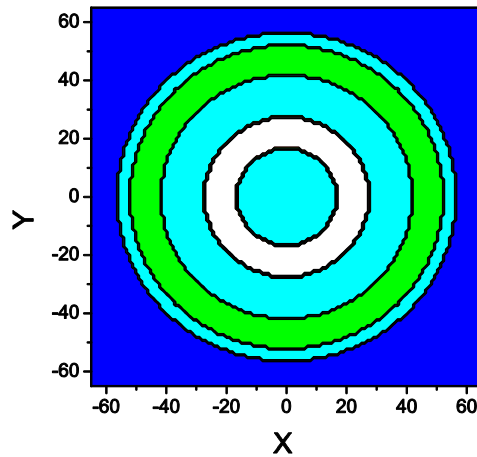
SF
MI ($\rho=1$)
MI ($\rho=2$)



	Number of atoms	V_0 (E_r)
a	7.1×10^5	15
b	8.9×10^5	15

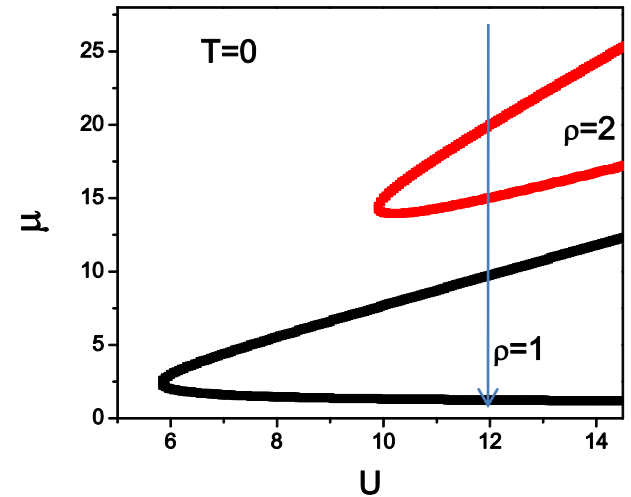
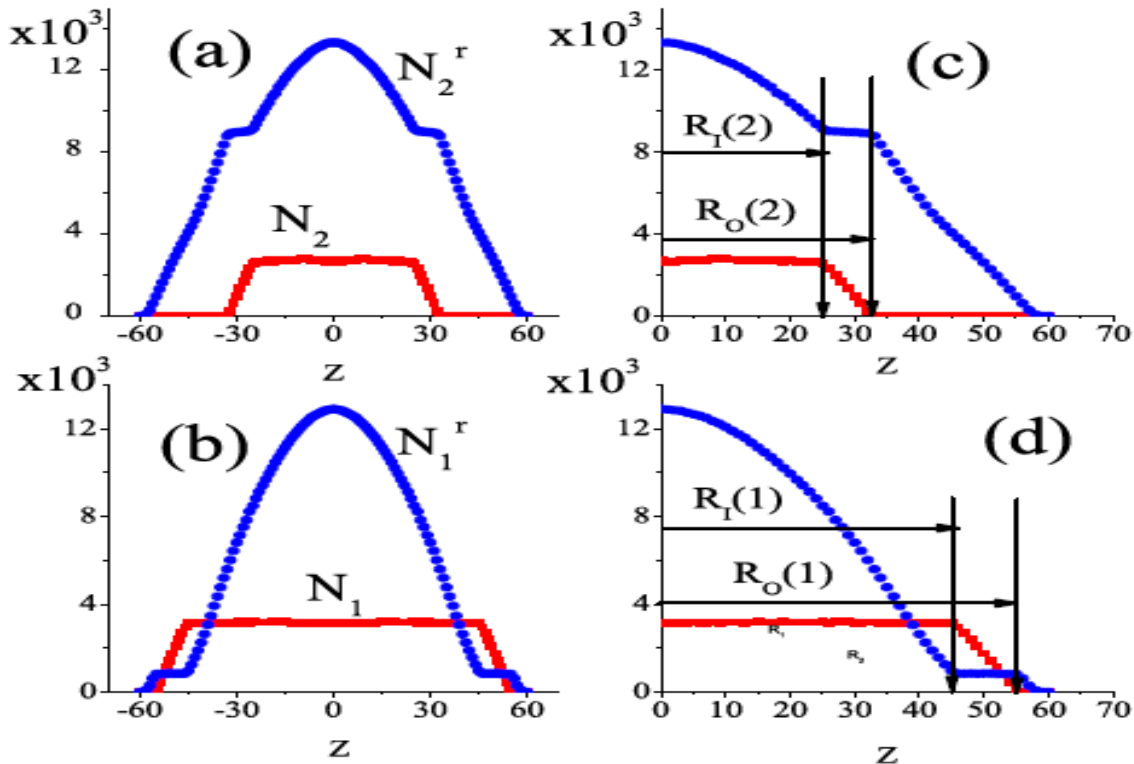
Phys. Rev. B **85**, 214524 (2012).

Bose-Hubbard model: Density Distribution & Shells: Result

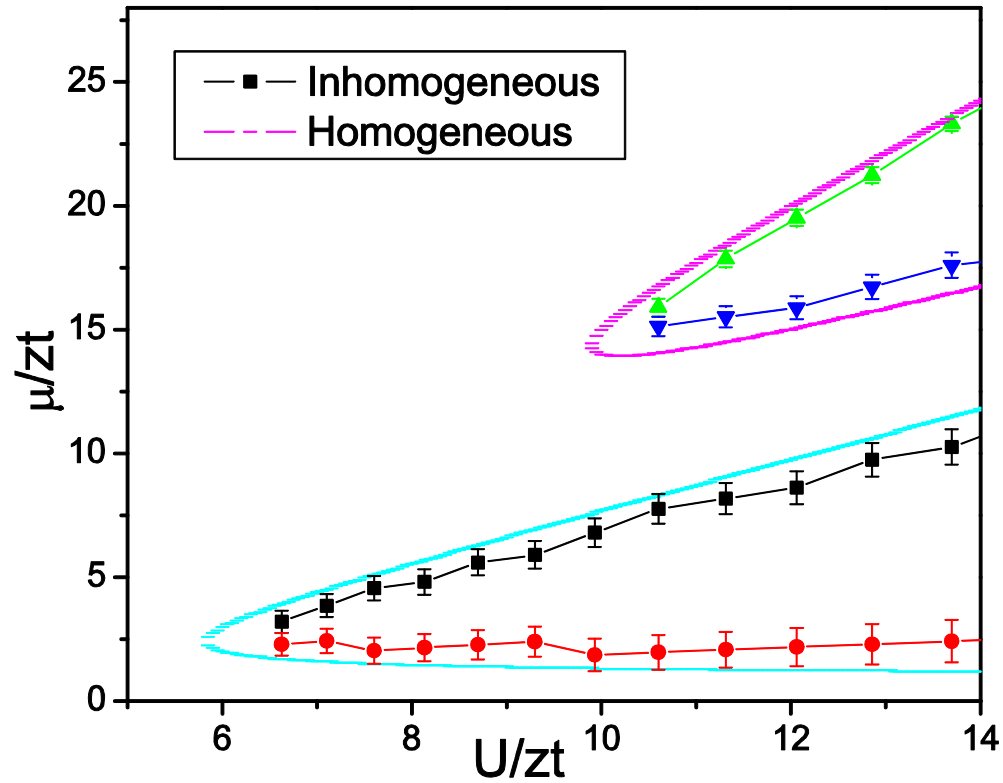


Phys. Rev. B **85**, 214524 (2012).

Bose-Hubbard model: Radii of MI shells



Bose-Hubbard model: Radii to phase diagram



Bose-Hubbard model: Excitations

$$H = \sum_i H_i^{MF} - t \sum_{\langle i, j \rangle} (\delta a_i^+ \delta a_j + h.c)$$

$$H_i^{MF} |i, \alpha \rangle = E_\alpha |i, \alpha \rangle$$

$$L_{\alpha\alpha'}^i = |i, \alpha \rangle \langle i, \alpha'| \quad \text{Define Projection Operator}$$

$$\hat{O}_i = \sum_{\alpha\alpha'} \langle i, \alpha | \hat{O}_i | i, \alpha' \rangle L_{\alpha\alpha'}^i \quad \text{Single Site Operator :}$$

$$H = \sum_{i\alpha} E_\alpha L_{\alpha\alpha}^i - \frac{t}{2} \sum_{\langle ij \rangle \alpha\alpha' \beta\beta'} T_{\alpha\alpha' \beta\beta'}^{ij} L_{\alpha\alpha'}^i L_{\beta\beta'}^j$$

$$T_{\alpha\alpha' \beta\beta'}^{ij} = \langle i, \alpha | \delta a_i^+ | i, \alpha' \rangle \langle j, \beta | \delta a_j | j, \beta' \rangle + h.c.$$

Bose-Hubbard model: Green's Function

$$G_{\alpha\alpha'\beta\beta'}^{ij}(t) = -i\theta(t) \langle [L_{\alpha\alpha'}^i(t), L_{\beta\beta'}^j(0)] \rangle$$

Equation of motion (RPA approximation):

$$(\omega - E_{\alpha} + E_{\alpha'})G_{\alpha\alpha'\beta\beta'}(q, \omega) + P_{\alpha\alpha'} \sum_{\mu\nu} T_{\alpha\alpha'\nu\mu}(q)G_{\mu\nu\beta\beta'}(q, \omega) = \frac{1}{2\pi} P_{\alpha\alpha'} \delta_{\alpha\beta'} \delta_{\alpha'\beta}$$

where $T_{\alpha\alpha'\beta\beta'}(q) = \varepsilon_q \left(T_{\alpha\alpha'\beta\beta'}^{ij} + T_{\beta\beta'\alpha\alpha'}^{ji} \right)$,

$$\varepsilon_q = -2t \sum_{j=x,y,z} \cos q_j$$

$$P_{\alpha\alpha'} = \langle L_{\alpha\alpha} \rangle - \langle L_{\alpha'\alpha'} \rangle$$

Bose-Hubbard model

Single Particle Green Functions

$$g_{i,j}(t) = -i\theta(t) \langle [a_i(t), a_j^+] \rangle$$

$$g(q, \omega) = \sum_{\alpha\alpha'\beta\beta'} \langle \alpha | a | \alpha' \rangle \langle \beta | a^+ | \beta' \rangle G_{\alpha\alpha'\beta\beta'}(q, \omega)$$

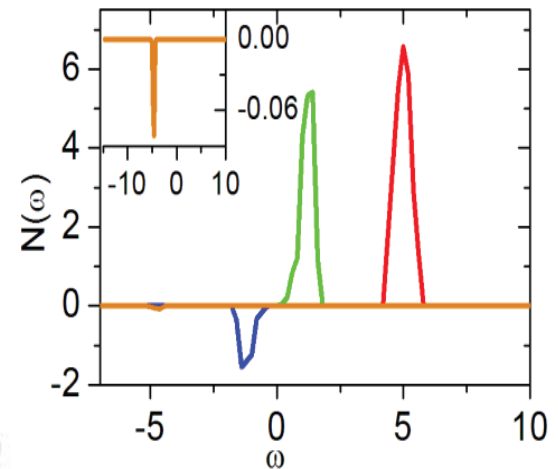
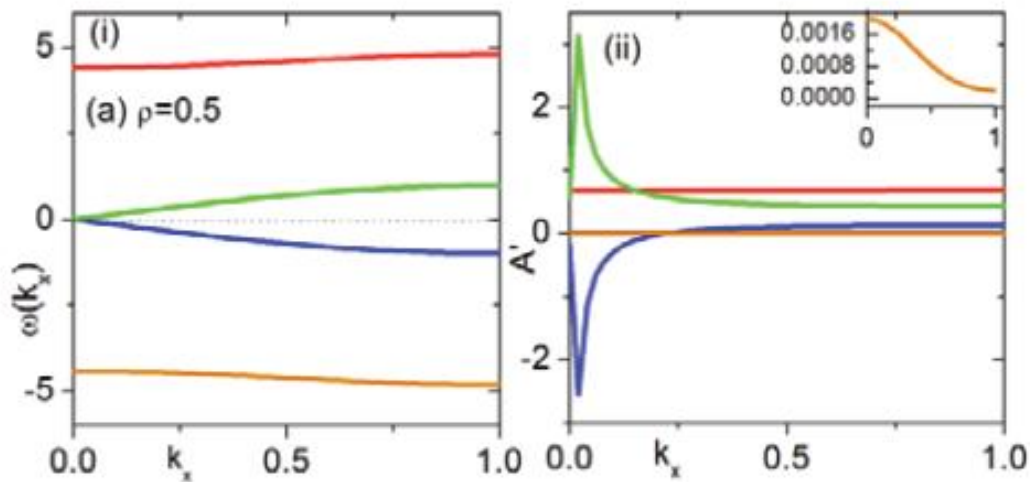
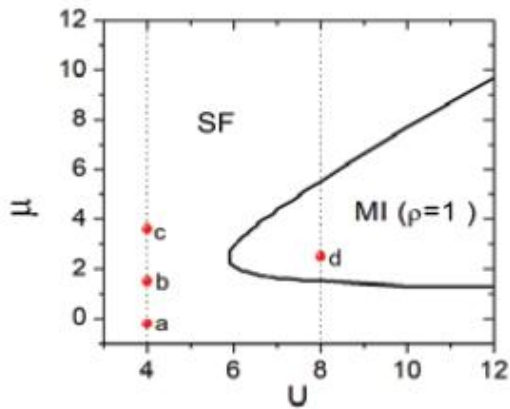
$$g(q, \omega) = \sum_r \frac{A(q)}{\omega - \omega_r(q)}$$

$$N(\omega) = -\frac{1}{\pi} \sum_q \text{Im } g(q, \omega^+)$$

rf-tunnelling current $I(\omega)$ spectroscopy is a powerful method to observe these excitation spectra.

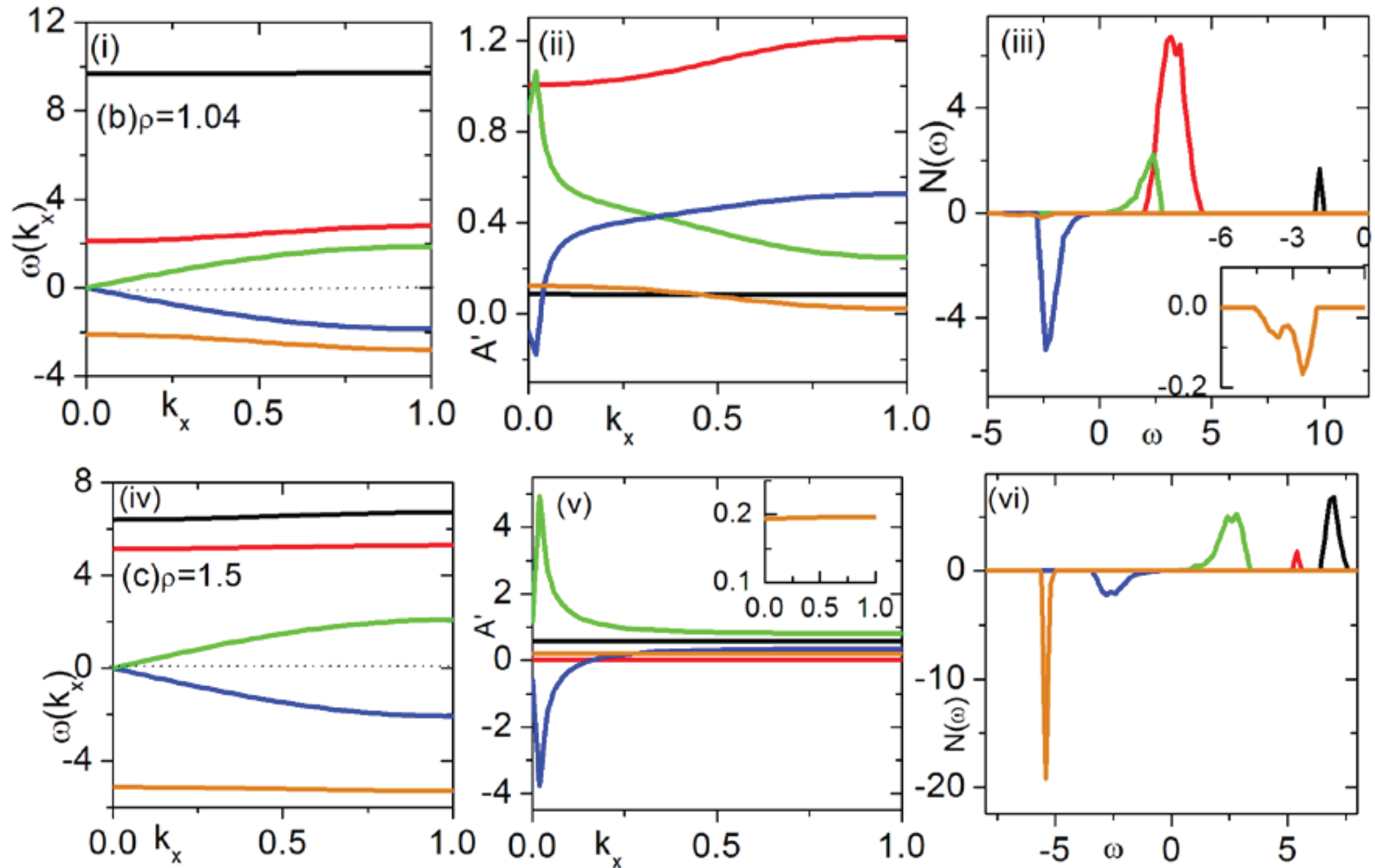
$$I(\omega) \propto N(\omega)$$

Bose-Hubbard model: Excitations – SF Phase



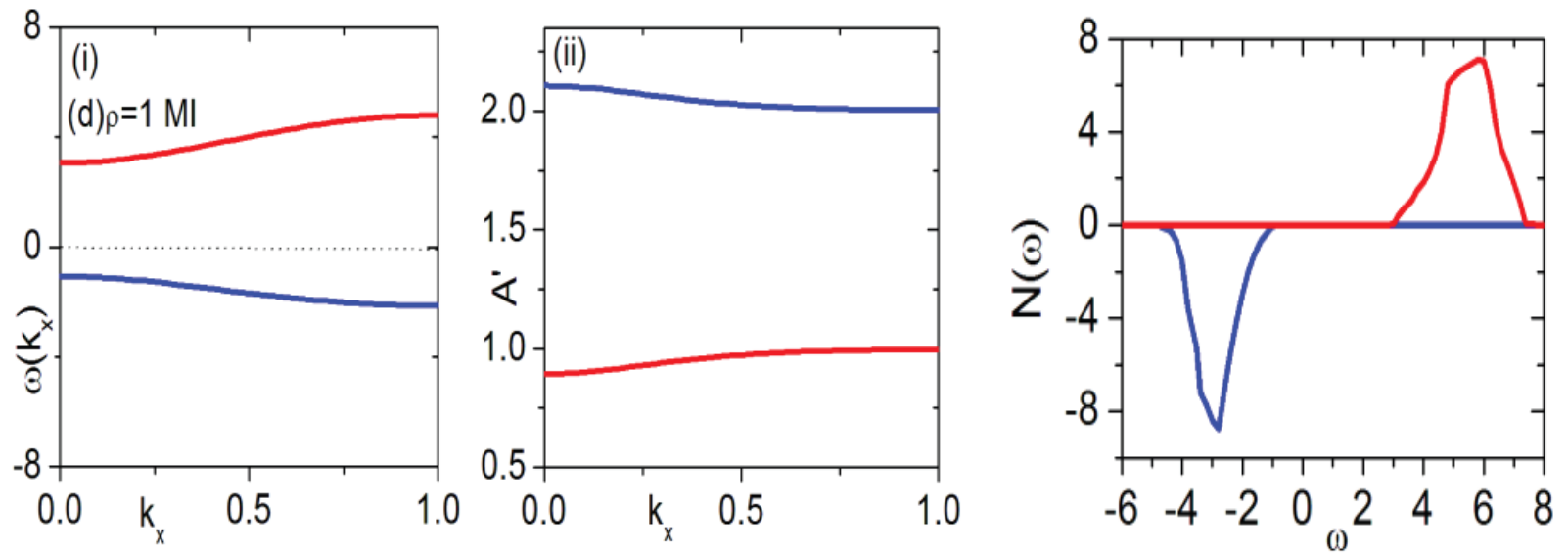
AIP Conference Proceedings **1832, 030009 (2017);**

Bose-Hubbard model: Excitations – SF Phase



AIP Conference Proceedings **1832, 030009 (2017);**

Bose-Hubbard model: Excitations - MI Phase



AIP Conference Proceedings **1832, 030009 (2017);**

Bose-Hubbard model: Finite Temperature

MFT

$$H_i^{MF} = \frac{U}{2} \hat{n}_i (\hat{n}_i - 1) - \mu n_i - (\psi_i^* a_i + \psi_i a_i^+) + |\psi_i|^2$$

$$H_i^{MF} |i, \alpha\rangle = E_\alpha |i, \alpha\rangle$$

Partition Function

$$Z = \sum_{\alpha} e^{-\beta E_\alpha}$$

Superfluid Order parameter $\psi = \sum_{\alpha} P_{\alpha} \langle \alpha | \hat{a} | \alpha \rangle$

Where $P_{\alpha} = \frac{e^{-\beta E_{\alpha}}}{Z}$

AIP Conference Proceedings (2018);

Bose-Hubbard model: Finite Temperature

RPA

$$H = \sum_i H_i^{MF} - t \sum_{\langle i,j \rangle} (\delta a_i^+ \delta a_j + h.c)$$

$$H_i^{MF} |i, \alpha \rangle = E_\alpha |i, \alpha \rangle$$

$$L_{\alpha\alpha'}^i = |i, \alpha \rangle \langle i, \alpha'| \quad \text{Define Projection Operator}$$

$$\hat{O}_i = \sum_{\alpha\alpha'} \langle i, \alpha | \hat{O}_i | i, \alpha' \rangle L_{\alpha\alpha'}^i \quad \text{Single Site Operator :}$$

$$H = \sum_{i\alpha} E_\alpha L_{\alpha\alpha}^i - \frac{t}{2} \sum_{\langle ij \rangle \alpha\alpha'\beta\beta'} T_{\alpha\alpha'\beta\beta'}^{ij} L_{\alpha\alpha'}^i L_{\beta\beta'}^j$$

$$T_{\alpha\alpha'\beta\beta'}^{ij} = \langle i, \alpha | \delta a_i^+ | i, \alpha' \rangle \langle j, \beta | \delta a_j | j, \beta' \rangle + h.c.$$

Bose-Hubbard model: Finite Temperature RPA

$$G_{\alpha\alpha'\beta\beta'}^{ij}(t) = -i\theta(t) \langle [L_{\alpha\alpha'}^i(t), L_{\beta\beta'}^j(0)] \rangle$$

Equation of motion (RPA approximation):

$$(\omega - E_{\alpha} + E_{\alpha'})G_{\alpha\alpha'\beta\beta'}(q, \omega) + P_{\alpha\alpha'} \sum_{\mu\nu} T_{\alpha\alpha'\nu\mu}(q)G_{\mu\nu\beta\beta'}(q, \omega) = \frac{1}{2\pi} P_{\alpha\alpha'} \delta_{\alpha\beta'} \delta_{\alpha'\beta}$$

where $T_{\alpha\alpha'\beta\beta'}(q) = \varepsilon_q \left(T_{\alpha\alpha'\beta\beta'}^{ij} + T_{\beta\beta'\alpha\alpha'}^{ji} \right),$

$$\varepsilon_q = -2t \sum_{j=x,y,z} \cos q_j$$

$$P_{\alpha\alpha'} = \langle L_{\alpha\alpha} \rangle - \langle L_{\alpha'\alpha'} \rangle$$

Bose-Hubbard model Finite Temperature RPA

$$G_{\alpha\alpha'\beta\beta'}(q, \omega) = \sum_r \frac{A(q)}{\omega - \omega_r(q)}$$

From this

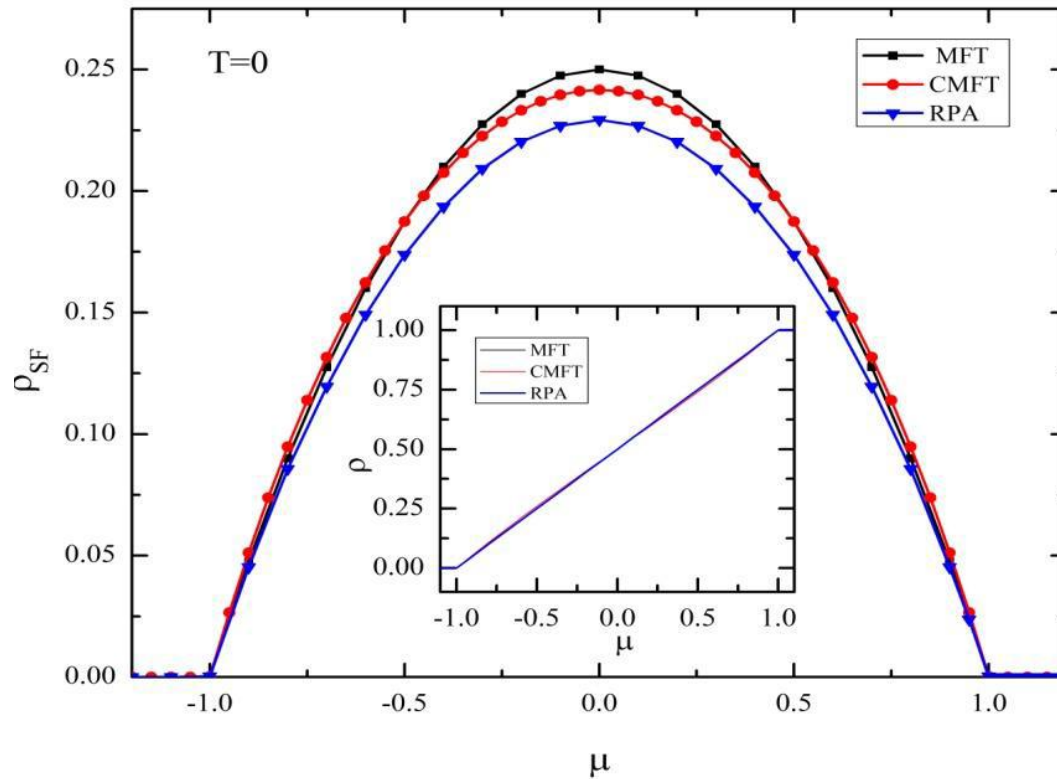
$$\psi = \sum P_\alpha \langle \alpha | \hat{a} | \alpha \rangle \quad \text{where}$$

$$P_\alpha = \frac{1}{(2\pi)^3} \sum_{r,q} A(q) f(\omega_r(q)) \quad \text{where}$$

$$f(\omega) = \left(e^{\beta\omega} - 1 \right)^{-1}$$

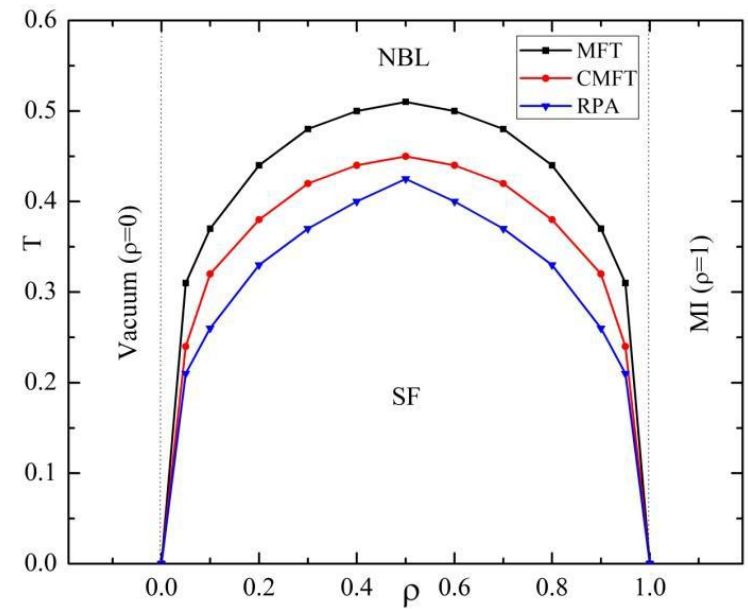
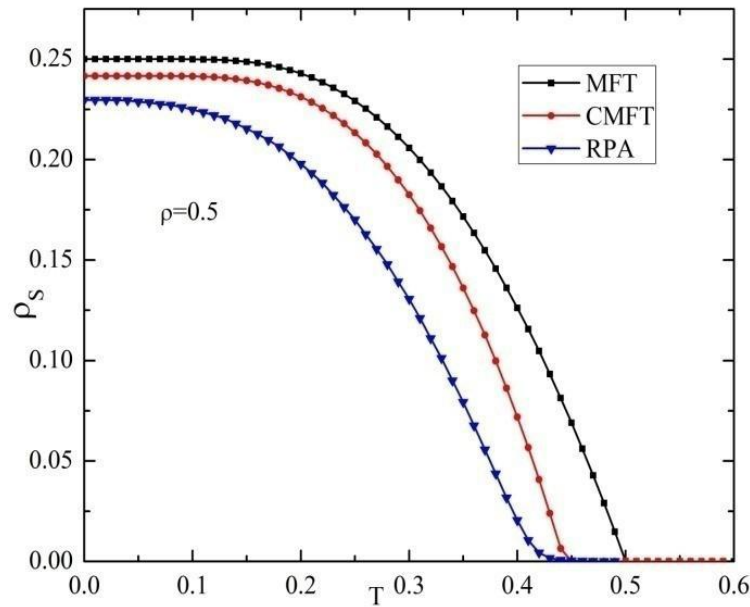
AIP Conference Proceedings (2018);

Hard Core Bose-Hubbard model: Finite Temperature Results



AIP Conference Proceedings (2018);

Hard Core Bose-Hubbard model: Finite Temperature Results



AIP Conference Proceedings (2018);

Bose-Hubbard model for two types of bosons

$$\begin{aligned} H = & -t_a \sum_{\langle i,j \rangle} (a_i^\dagger a_j + h.c.) + \frac{U_a}{2} \sum_i n_i^a (n_i^a - 1) \\ & -t_b \sum_{\langle i,j \rangle} (b_i^\dagger b_j + h.c.) + \frac{U_b}{2} \sum_i n_i^b (n_i^b - 1) \\ & + U_{ab} \sum_i n_i^a n_i^b - \mu_a \sum_i n_i^a - \mu_b \sum_i n_i^b \end{aligned}$$

Bose-Hubbard model for two types of bosons

Mean field theory

$$a_i^+ a_j \approx \langle a_i^+ \rangle a_j + a_i^+ \langle a_j \rangle - \langle a_i^+ \rangle \langle a_j \rangle$$

$$b_i^+ b_j \approx \langle b_i^+ \rangle b_j + b_i^+ \langle b_j \rangle - \langle b_i^+ \rangle \langle b_j \rangle$$

Superfluid order parameters

$$\psi_i^a = \langle a_i \rangle \quad \psi_i^b = \langle b_i \rangle$$

$$H = \sum_i H_i^{MF} - t_a \sum_{\langle i,j \rangle} (\delta a_i^+ \delta a_j + h.c) - t_b \sum_{\langle i,j \rangle} (\delta b_i^+ \delta b_j + h.c)$$

Bose-Hubbard model for two types of bosons

Mean field theory

$$\begin{aligned}
 H_i^{MF} = & \frac{U_a}{2} \hat{n}_i^a (\hat{n}_i^a - 1) - \mu_a n_i^a - (\psi_i^{a*} a_i + \psi_i^a a_i^+) + |\psi_i^a|^2 \\
 & + \frac{U_b}{2} \hat{n}_i^b (\hat{n}_i^b - 1) - \mu_b n_i^b - (\psi_i^{b*} b_i + \psi_i^b b_i^+) + |\psi_i^b|^2 \\
 & + U_{ab} \hat{n}_i^a \hat{n}_i^b
 \end{aligned}$$

- Superfluid density $(\rho_s^a)_i = |\psi_i^a|^2$ $(\rho_s^b)_i = |\psi_i^b|^2$
- Density $\rho_i^a = \langle \hat{n}_i^a \rangle$ $\rho_i^b = \langle \hat{n}_i^b \rangle$

Bose-Hubbard model for two types of bosons

Mean field theory

$$\begin{aligned}
 H_i^{MF} = & \frac{U_a}{2} \hat{n}_i^a (\hat{n}_i^a - 1) - \mu_a n_i^a - (\psi_i^{a*} a_i + \psi_i^a a_i^+) + |\psi_i^a|^2 \\
 & + \frac{U_b}{2} \hat{n}_i^b (\hat{n}_i^b - 1) - \mu_b n_i^b - (\psi_i^{b*} b_i + \psi_i^b b_i^+) + |\psi_i^b|^2 \\
 & + U_{ab} \hat{n}_i^a \hat{n}_i^b
 \end{aligned}$$

- Superfluid density $(\rho_s^a)_i = |\psi_i^a|^2$ $(\rho_s^b)_i = |\psi_i^b|^2$
- Density $\rho_i^a = \langle \hat{n}_i^a \rangle$ $\rho_i^b = \langle \hat{n}_i^b \rangle$

Bose-Hubbard model for two types of bosons

Mean field theory: Phases

$$SF_A + SF_B \quad \rho_S^a > 0, \quad \rho_S^b > 0$$

$$MI_A + MI_B \quad \rho_S^a = 0, \quad \rho_S^b = 0$$

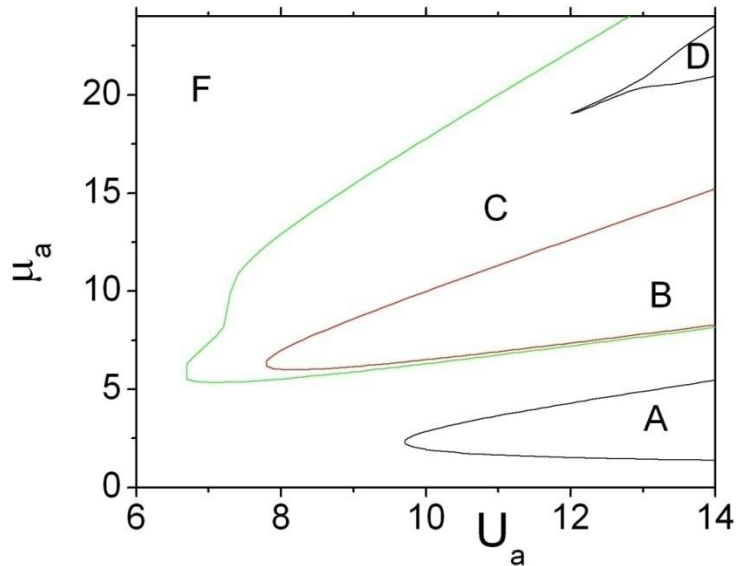
$$MI_A + SF_B \quad \rho_S^a = 0, \quad \rho_S^b > 0$$

$$MI_B + SF_A \quad \rho_S^b = 0, \quad \rho_S^a > 0$$

Bose-Hubbard model for two types of bosons

Mean field theory: Phases

- A : MI-1 $\rho_A + \rho_B = 1$
- B : MI_A, MI_B $\rho_A = 1, \rho_B = 1$
- C : MI_A, SF_B $\rho_A = 1, \rho_B > 1$
- D : MI_A, MI_B $\rho_A = 1, \rho_B = 2$
- F : SF_A, SF_B



Phase diagram, and plots versus μ of the for
 $U_{ab} = 0.5U_a$, $U_b = 0.75U_a$, $\mu_a = \mu_b = \mu$.

Bose-Hubbard model for two types of bosons

Excitations

$$H = \sum_i H_i^{MF} - t_a \sum_{\langle i,j \rangle} (\delta a_i^+ \delta a_j + h.c) \\ - t_b \sum_{\langle i,j \rangle} (\delta b_i^+ \delta b_j + h.c)$$

$$H_i^{MF} |i, \alpha \rangle = E_i |i, \alpha \rangle$$

Define Projection Operator $L_{\alpha\alpha'}^i = |i, \alpha \rangle \langle i, \alpha'|$

Single Site Operator : $\hat{O}_i = \sum_{\alpha\alpha'} \langle i, \alpha | \hat{O}_i | i, \alpha' \rangle L_{\alpha\alpha'}^i$

Bose-Hubbard model for two types of bosons

Excitations

$$H = \sum_{i\alpha} E_{\alpha} L_{\alpha\alpha}^i - \frac{1}{2} \sum_{\langle ij \rangle \alpha\alpha' \beta\beta'} T_{\alpha\alpha' \beta\beta'}^{ij} L_{\alpha\alpha'}^i L_{\beta\beta'}^j$$

$$T_{\alpha\alpha' \beta\beta'}^{ij} = \langle i, \alpha | \delta a_i^+ | i, \alpha' \rangle \langle j, \beta | \delta a_j | j, \beta' \rangle + h.c.$$

$$+ \langle i, \alpha | \delta b_i^+ | i, \alpha' \rangle \langle j, \beta | \delta b_j | j, \beta' \rangle + h.c.$$

$$\delta a = a - \langle a \rangle$$

$$\delta b = b - \langle b \rangle$$

Bose-Hubbard model for two types of bosons

Green Functions + RPA

$$G_{\alpha\alpha'\beta\beta'}^{ij}(t) = -i\theta(t) \langle [L_{\alpha\alpha'}^i(t), L_{\beta\beta'}^j(0)] \rangle$$

Equation of motion (RPA approximation):

$$(\omega - E_\alpha + E_{\alpha'})G_{\alpha\alpha'\beta\beta'}(q, \omega) + P_{\alpha\alpha'} \sum_{\mu\nu} T_{\alpha\alpha'\nu\mu}(q)G_{\mu\nu\beta\beta'}(q, \omega) = \frac{1}{2\pi} P_{\alpha\alpha'} \delta_{\alpha\beta'} \delta_{\alpha'\beta}$$

where $T_{\alpha\alpha'\beta\beta'}(q) = \varepsilon_q \left(T_{\alpha\alpha'\beta\beta'}^{ij} + T_{\beta\beta'\alpha\alpha'}^{ji} \right)$,

$$\varepsilon_q = -2t \sum_{j=x,y,z} \cos q_j$$

$$P_{\alpha\alpha'} = \langle L_{\alpha\alpha} \rangle - \langle L_{\alpha'\alpha'} \rangle$$

Bose-Hubbard model for two types of bosons

Single Particle Green Functions

$$g_{i,j}^{aa}(t) = -i\theta(t) \langle [a_i(t), a_j^+] \rangle$$

$$g_{i,j}^{bb}(t) = -i\theta(t) \langle [b_i(t), b_j^+] \rangle$$

$$g^{aa}(q, \omega) = \sum_{\alpha\alpha'\beta\beta'} \langle \alpha | a | \alpha' \rangle \langle \beta | a^+ | \beta' \rangle G_{\alpha\alpha'\beta\beta'}(q, \omega)$$

$$g^{bb}(q, \omega) = \sum_{\alpha\alpha'\beta\beta'} \langle \alpha | b | \alpha' \rangle \langle \beta | b^+ | \beta' \rangle G_{\alpha\alpha'\beta\beta'}(q, \omega)$$

Bose-Hubbard model for two types of bosons

Excitations, Spectral Weights + DOS

Solving the equation of motion for the Green function and writing it in the form

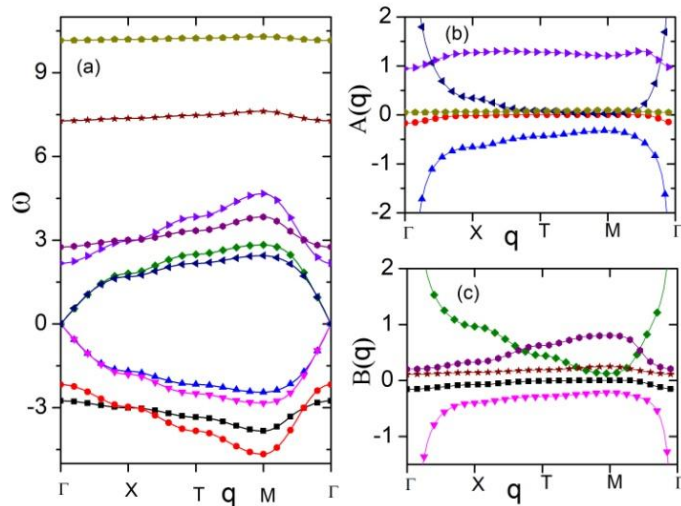
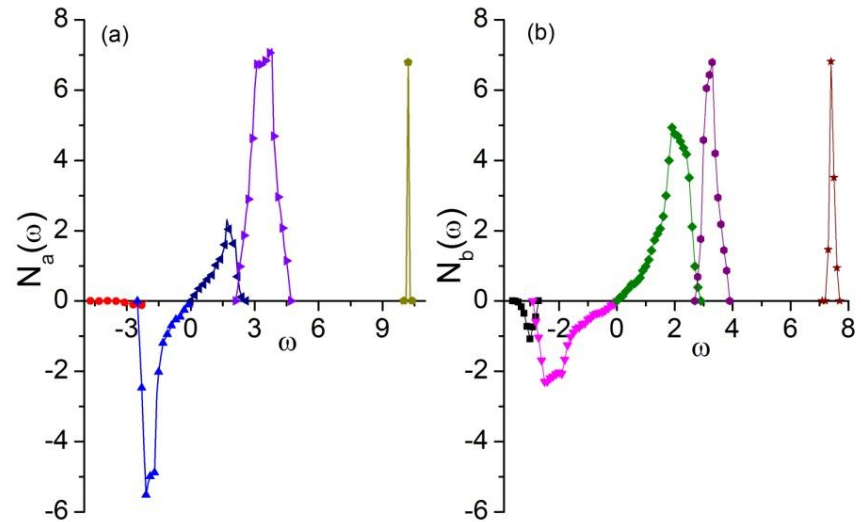
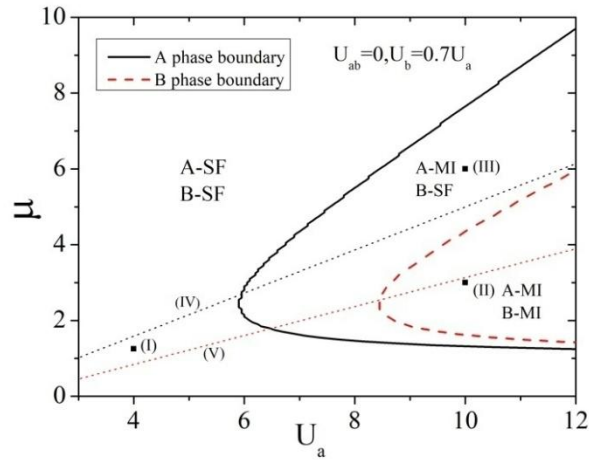
$$g^{aa}(q, \omega) = \sum_r \frac{A(q)}{\omega - \omega_r(q)}$$

$$g^{bb}(q, \omega) = \sum_r \frac{B(q)}{\omega - \omega_r(q)}$$

$$N_a(\omega) = -\frac{1}{\pi} \sum_q \text{Im} g^{aa}(q, \omega^+)$$

Bose-Hubbard model for two types of bosons

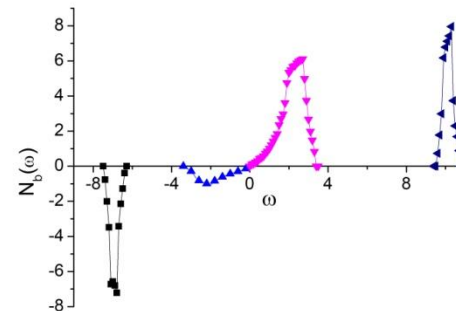
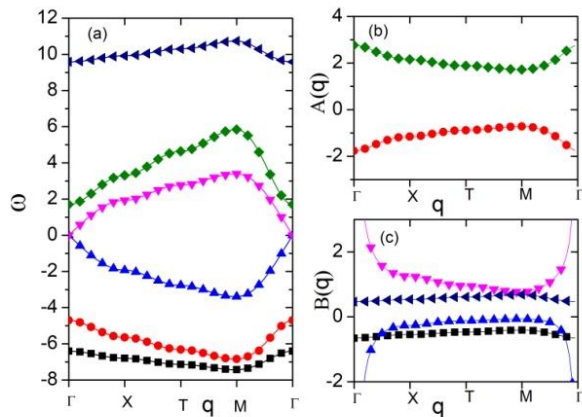
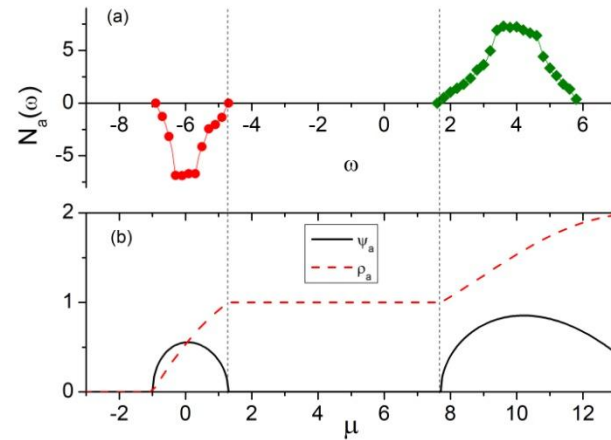
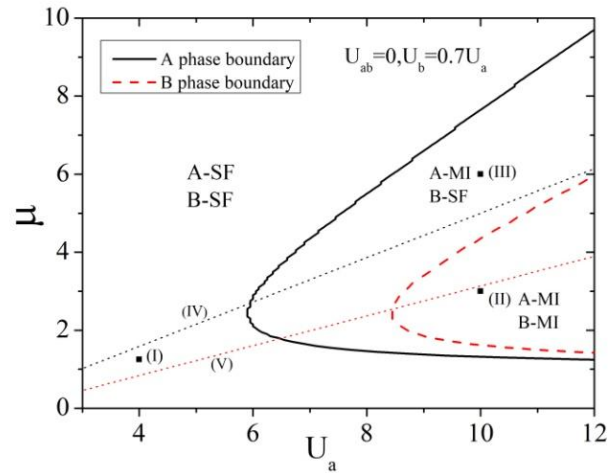
Excitations



SF_A, SF_B

Bose-Hubbard model for two types of bosons

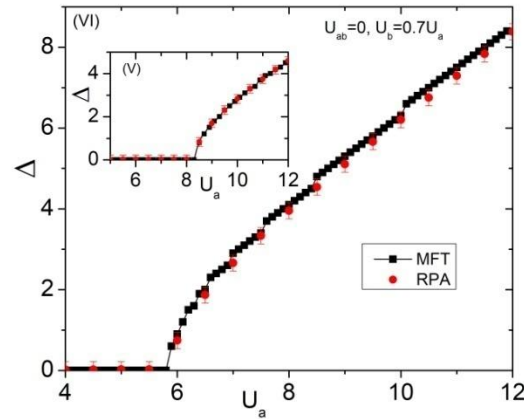
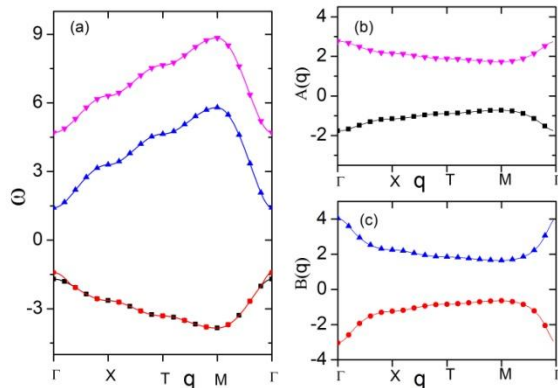
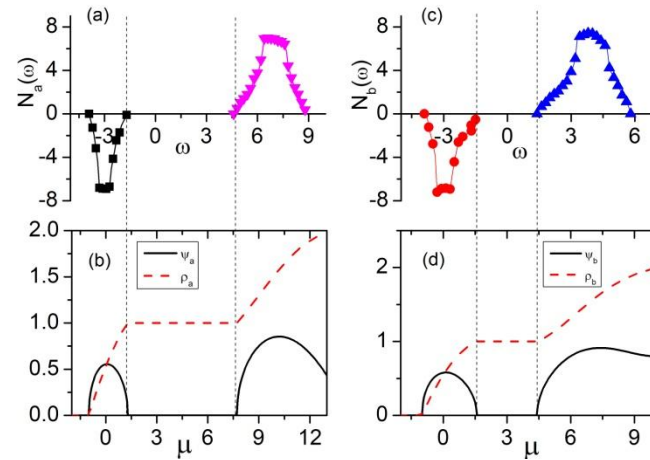
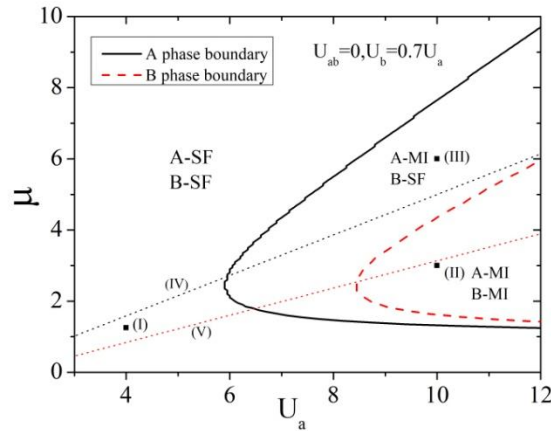
Excitations



MI_A, SF_B

Bose-Hubbard model for two types of bosons

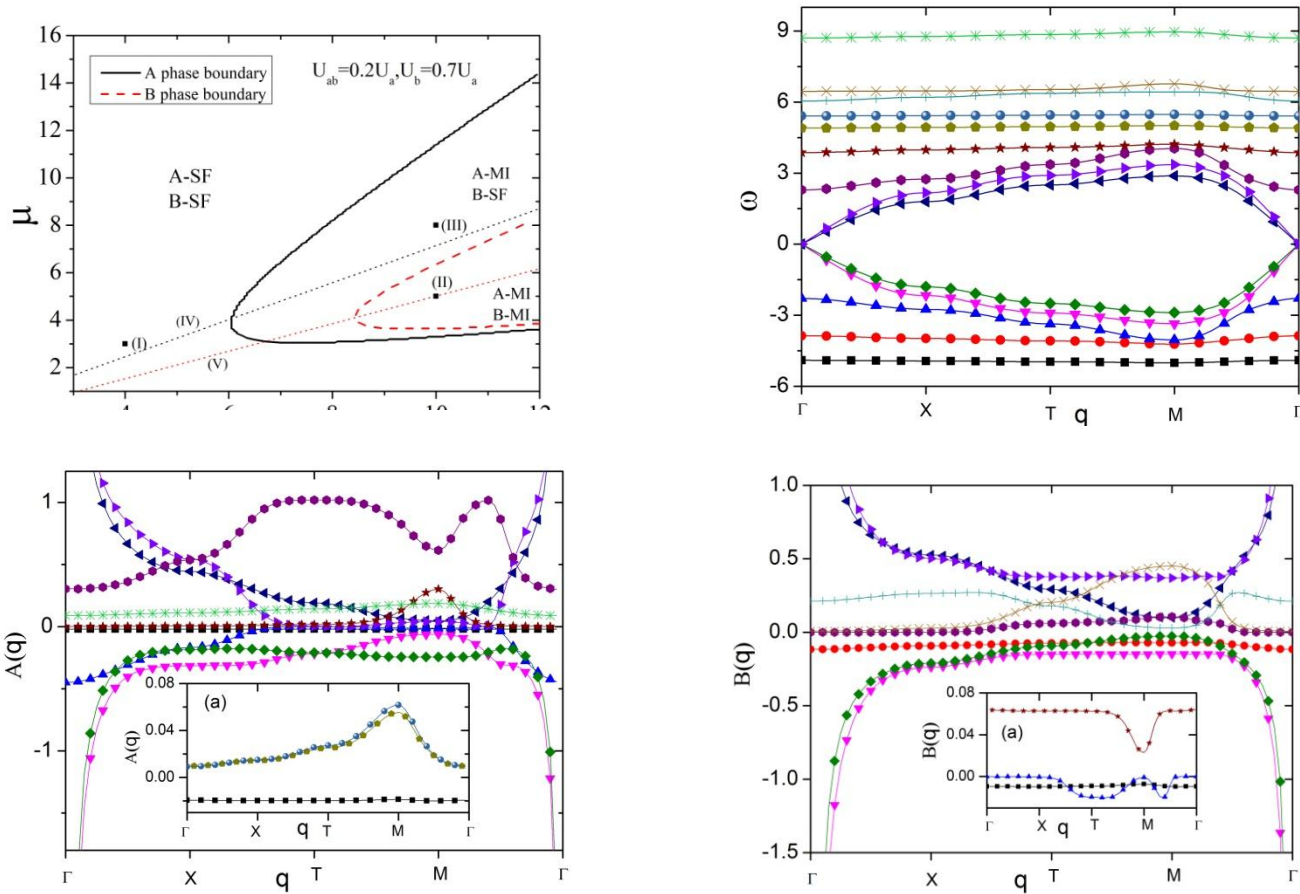
Excitations



MI_A, MI_B

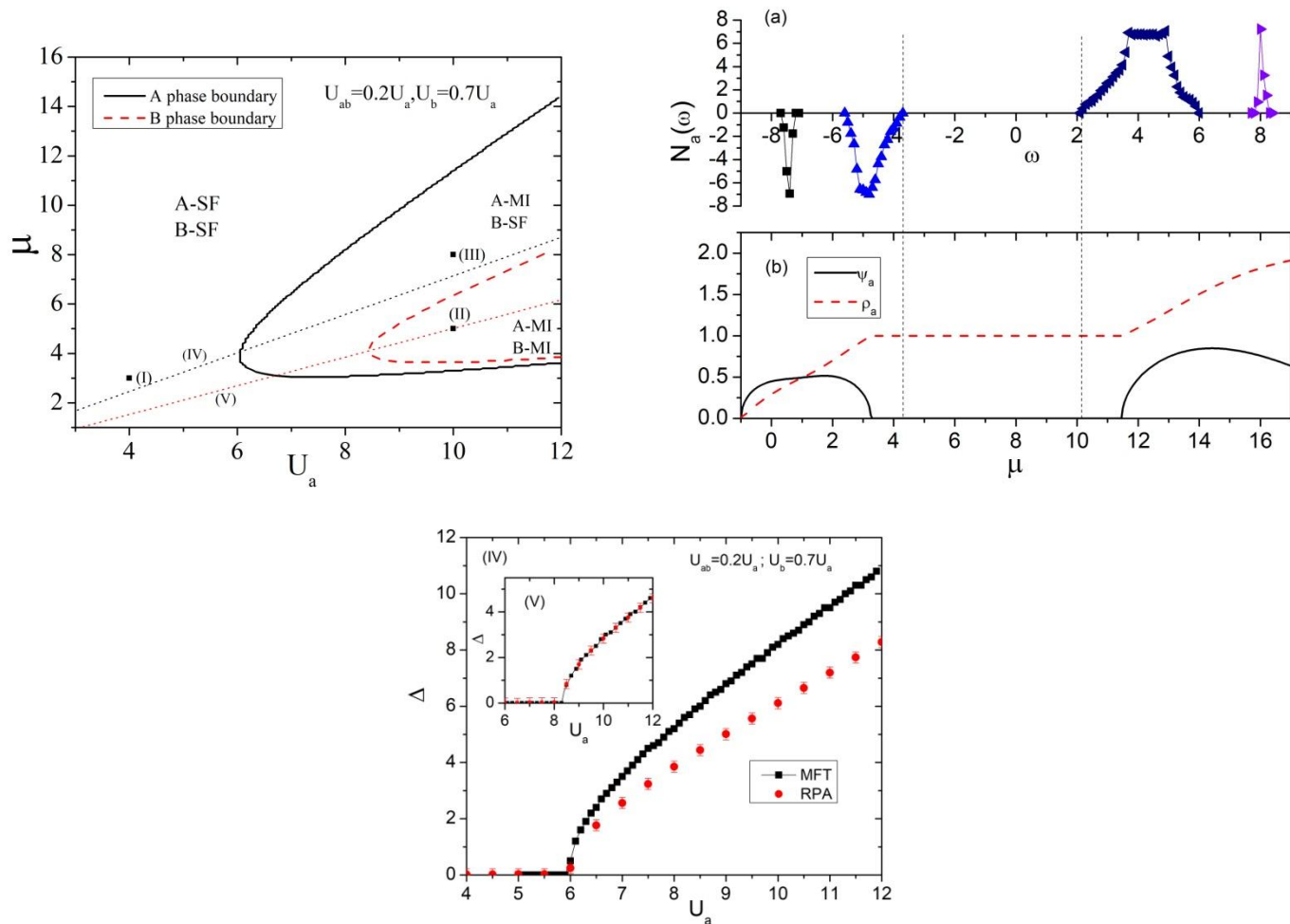
Bose-Hubbard model for two types of bosons

Excitations



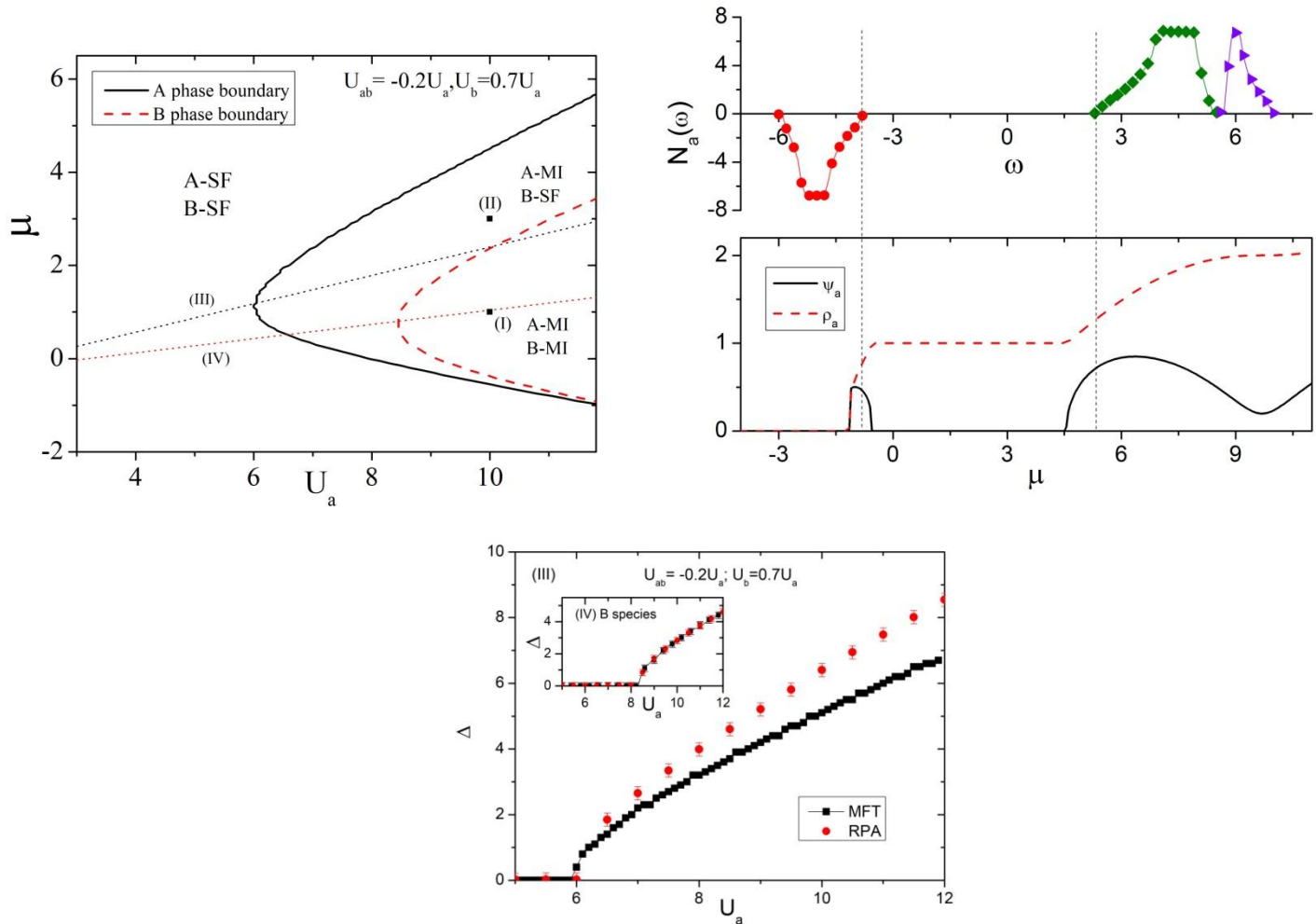
Bose-Hubbard model for two types of bosons

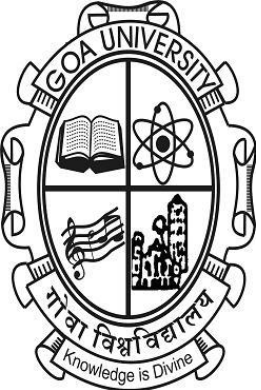
Excitations



Bose-Hubbard model for two types of bosons

Excitations



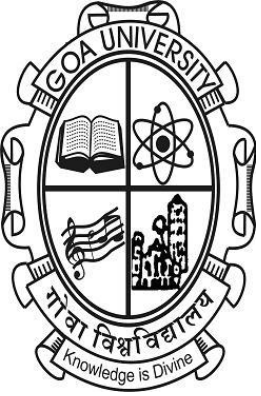


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Ramesh V Pai

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