



An almost convincing scheme for discriminating the preparation basis of quantum ensemble and why it will not work

Sandeep K. Goyal

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IISER Mohali

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# Quantum states

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$$|\psi\rangle$$



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$$|\psi\rangle = \alpha |\uparrow\rangle + \beta |\downarrow\rangle$$

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$$\rho = \sum_n p_n |\psi_n\rangle \langle\psi_n|, \quad \sum_n p_n = 1, \quad p_n \geq 0$$



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$$\text{tr}\rho = 1$$



# Density operators

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$$\rho = \sum_n p_n |\psi_n\rangle \langle \psi_n|$$

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$$\sqrt{p_n} |\psi_n\rangle = \sum_m W_{nm} \sqrt{q_m} |\phi_m\rangle$$



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$$\sum_n W_{nm} W_{nm'}^* = \delta_{mm'} \equiv W^\dagger W = \mathbb{1}$$



# Quantum Measurement

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- Consider the state  $|\psi\rangle = \alpha |\uparrow\rangle + \beta |\downarrow\rangle$
- Upon measurement in basis  $\{|\uparrow\rangle, |\downarrow\rangle\}$ , the state will collapse to  $|\uparrow\rangle$  with probability  $|\alpha|^2$  and to state  $|\downarrow\rangle$  with probability  $|\beta|^2$ .



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- After the measurement the quantum system lose all the information about the state  $|\psi\rangle$ .



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- A more general measurement approach is to make an ancillary system interact with the quantum system under observation and perform measurement on the ancillary system afterward



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- A more general measurement approach is to make an ancillary system interact with the quantum system under observation and perform measurement on the ancillary system afterward
- For example, we start with a system in state  $|\psi\rangle = \alpha |0\rangle + \beta |1\rangle$  and an ancillary system in state  $|0\rangle$ .



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- A more general measurement approach is to make an ancillary system interact with the quantum system under observation and perform measurement on the ancillary system afterward
- For example, we start with a system in state  $|\psi\rangle = \alpha |0\rangle + \beta |1\rangle$  and an ancillary system in state  $|0\rangle$ .
- The joint state of the system and ancilla is  $|\psi\rangle |0\rangle$ .



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If  $U$  is the unitary operator that characterizes the interaction between the system and ancilla then the state after the interaction reads



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If  $U$  is the unitary operator that characterizes the interaction between the system and ancilla then the state after the interaction reads

$$U(|\psi\rangle|0\rangle) = \left(\sqrt{p_0}\alpha|0\rangle + \sqrt{1-p_0}\beta|1\rangle\right)|0\rangle + \left(\sqrt{1-p_0}\alpha|0\rangle + \sqrt{p_0}\beta|1\rangle\right)|1\rangle.$$



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Upon measurement the ancillary system will collapse to  $\{|0\rangle, |1\rangle\}$  which result in the state of the system





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Upon measurement the ancillary system will collapse to  $\{|0\rangle, |1\rangle\}$  which result in the state of the system

$$|\psi\rangle_+ = \sqrt{p_0}\alpha |0\rangle + \sqrt{1-p_0}\beta |1\rangle \equiv M_+ |\psi\rangle,$$

$$|\psi\rangle_- = \sqrt{1-p_0}\alpha |0\rangle + \sqrt{p_0}\beta |1\rangle \equiv M_- |\psi\rangle.$$



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$$M_+ = \begin{pmatrix} \sqrt{p_0} & 0 \\ 0 & \sqrt{1-p_0} \end{pmatrix}, \quad M_- = \begin{pmatrix} \sqrt{1-p_0} & 0 \\ 0 & \sqrt{p_0} \end{pmatrix}.$$



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$$|\psi\rangle \rightarrow \begin{cases} M_+ |\psi\rangle & \text{with } p_+ = \langle\psi| M_+^\dagger M_+ |\psi\rangle \\ M_- |\psi\rangle & \text{with } p_- = \langle\psi| M_-^\dagger M_- |\psi\rangle \end{cases}$$



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- Since,  $p_+ + p_- = 1$  for all the states  $|\psi\rangle$  we have  $M_+^\dagger M_+ + M_-^\dagger M_- = \mathbb{1}$



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- Since,  $p_+ + p_- = 1$  for all the states  $|\psi\rangle$  we have  $M_+^\dagger M_+ + M_-^\dagger M_- = \mathbb{1}$
- The expectation value of the operator  $\sigma_z$  is proportional to the probabilities  $p_+$  and  $p_-$ , i.e.,

$$\langle\sigma_z\rangle \propto p_+ - p_-$$



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$$M_+ = \begin{pmatrix} \sqrt{p_0} & 0 \\ 0 & \sqrt{1-p_0} \end{pmatrix}, \quad M_- = \begin{pmatrix} \sqrt{1-p_0} & 0 \\ 0 & \sqrt{p_0} \end{pmatrix}$$



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$$|\psi\rangle \rightarrow \begin{cases} M_+ |\psi\rangle \rightarrow \begin{cases} M_+^2 |\psi\rangle \\ M_- M_+ |\psi\rangle \end{cases} \\ M_- |\psi\rangle \rightarrow \begin{cases} M_+ M_- |\psi\rangle \\ M_-^2 |\psi\rangle \end{cases} \end{cases}$$





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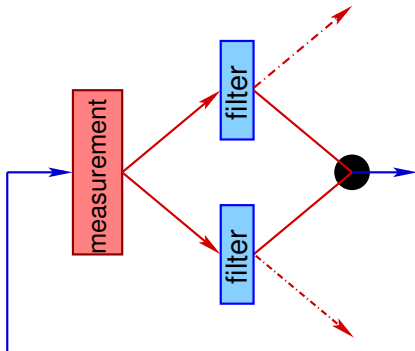
$$M_+ M_- = M_- M_+ = \sqrt{p_0(1-p_0)} \mathbb{1}$$



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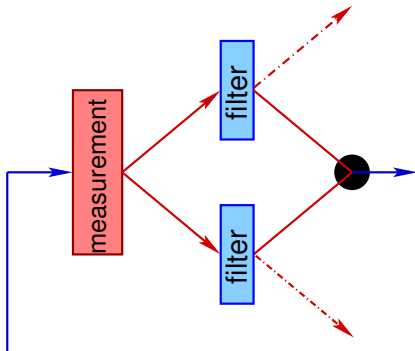
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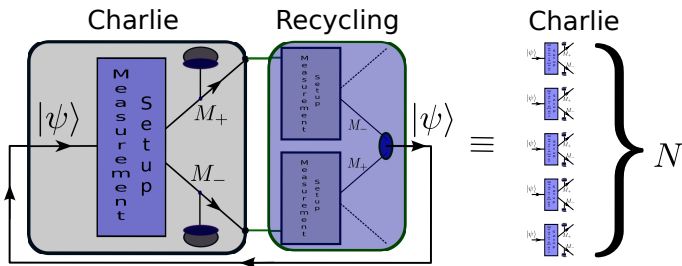


$$p_{rev} = 2p_0(1 - p_0)$$

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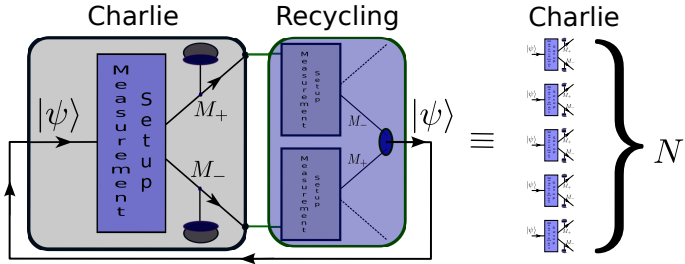
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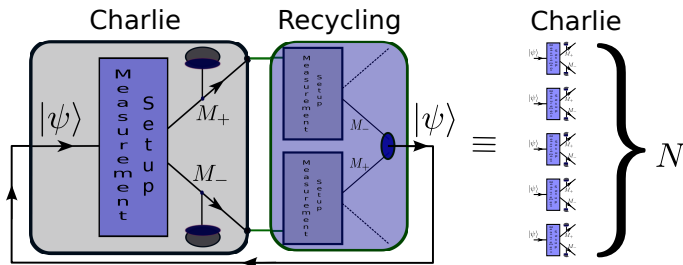


$$N_+ = \#M_+, \quad N_- = \#M_-, \quad N_+ + N_- = N$$

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$$N_+ = \#M_+, \quad N_- = \#M_-, \quad N_+ + N_- = N$$

$$n = (N_+ - N_-)/N$$



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If  $N \rightarrow \infty$



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$$|\psi\rangle = |0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \Rightarrow n = 1 - 2p_0$$





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$$|\psi\rangle = |\pm\rangle = \frac{1}{\sqrt{2}}(|0\rangle \pm |1\rangle) \Rightarrow n = 0$$



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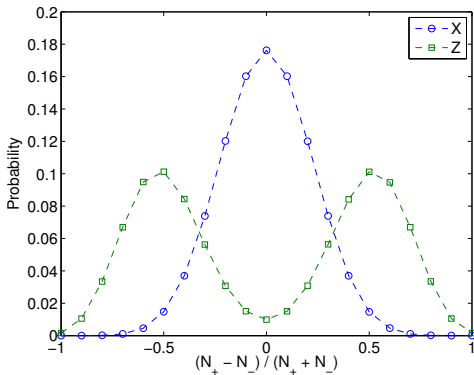
For finite  $N$  the value of  $n$  will fall on a Gaussian curve centered around the expected value of  $n$  with a width proportional to  $1/\sqrt{N}$



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# Estimating preparation basis

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$$\rho = \frac{1}{2} \mathbb{1}_2$$

$$\rho = \frac{1}{2}(|0\rangle\langle 0| + |1\rangle\langle 1|)$$



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$$\rho = \frac{1}{2}(|+\rangle\langle +| + |-\rangle\langle -|)$$

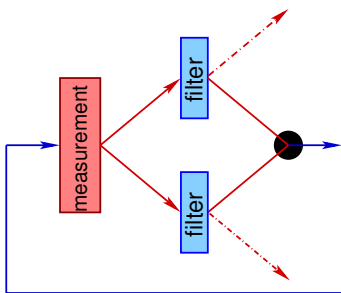




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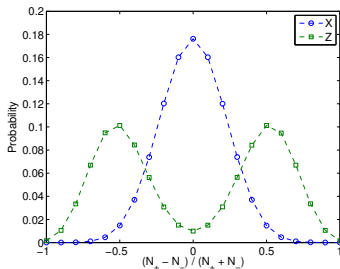
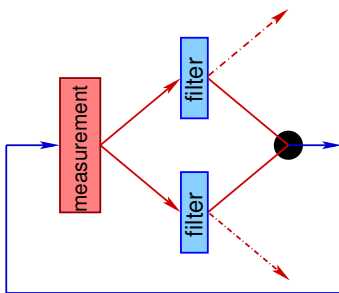
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# Superluminal communication

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A method to estimate the preparation basis can result in superluminal communication.



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A method to estimate the preparation basis can result in superluminal communication.

- Alice wants to send one bit of information  $\{0, 1\}$  to Bob without using classical communication



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A method to estimate the preparation basis can result in superluminal communication.

- Alice wants to send one bit of information  $\{0, 1\}$  to Bob without using classical communication
- They have infinite supply of maximally entangled states  $|\Phi\rangle$

$$|\Phi\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$$



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- If Alice wants to send '0' to Bob, she performs measurement in the eigenbasis of  $\sigma_z$ , i.e.,  $\{|0\rangle, |1\rangle\}$



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$$|\Phi\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$$

- If Alice wants to send '0' to Bob, she performs measurement in the eigenbasis of  $\sigma_z$ , i.e.,  $\{|0\rangle, |1\rangle\}$
- If Alice wants to send '1' to Bob, she performs measurement in the eigenbasis of  $\sigma_x$ , i.e.,  $\{|+\rangle, |-\rangle\}$



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- As soon as Alice performs measurement the state of the qubits in the Bob's possession acquire the state  $\mathbb{1}/2$  which is either

$$\rho = \frac{1}{2}(|0\rangle\langle 0| + |1\rangle\langle 1|)$$





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# Where is the problem?

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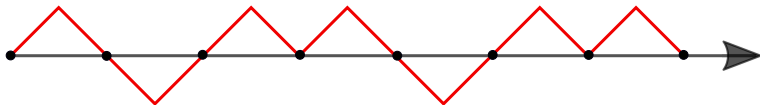
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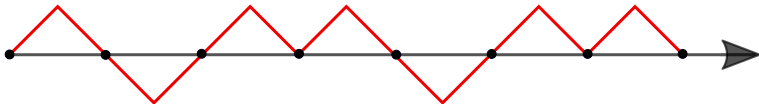




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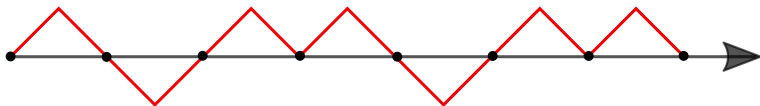
- All the trajectories with  $N$  retrievals have the same probability  $p_0^N (1 - p_0)^N$



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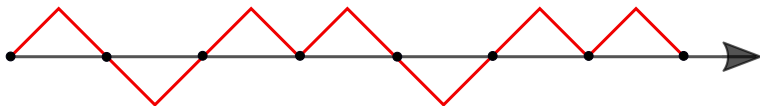
- All the trajectories with  $N$  retrievals have the same probability  $p_0^N (1 - p_0)^N$
- The number of possible trajectories when  $N_+$  number of + clicks occurs and  $N_- = N - N_+$  number of - clicks, is  $N! / (N_+! (N - N_+)!)$



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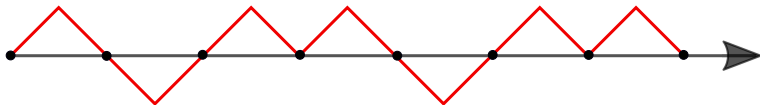
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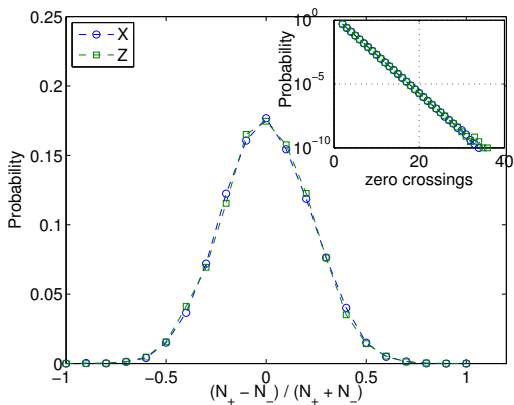
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- The number of possible trajectories when  $N_+$  number of + clicks occurs and  $N_- = N - N_+$  number of - clicks, is  $N! / (N_+! (N - N_+)!)$
- This number is largest when  $N_+ = N_-$
- Thus, the trajectories which have  $N_+ = N_-$  dominates the statistics



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- Charlie is tasked with the measurement of a quantum system in an unknown state  $|\psi\rangle$



# The paradox

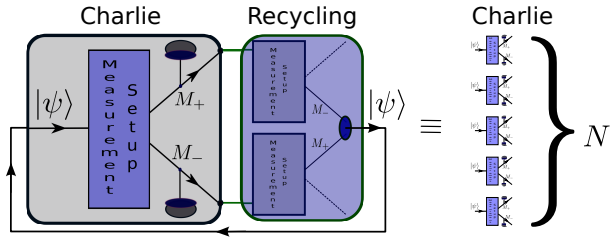
An almost convincing scheme for discriminating the preparation basis of quantum ensemble and why it will not work

Sandeep  
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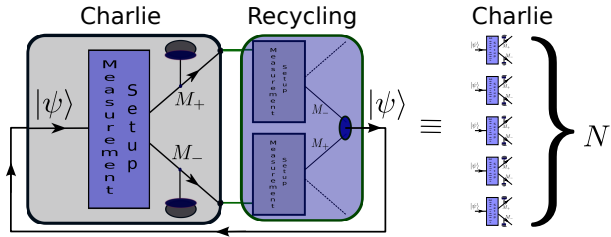


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# The paradox

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- However, the two ways of preparing the states result in entirely different outcomes

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The paradox is more interesting when Charlie know the state  $|\psi\rangle$ . In that case he will know the means used to prepare the state just by looking at the statistics.



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The paradox is more interesting when Charlie know the state  $|\psi\rangle$ . In that case he will know the means used to prepare the state just by looking at the statistics.

Even though the systems and the states were identical in both the cases.



# Reference

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“How measurement reversal could erroneously suggest the capability to discriminate the preparation basis of a quantum ensemble”  
Sandeep K. Goyal, Rajeev Singh, and Sibasish Ghosh  
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