

TWISTING REALITY AND FODO

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GEOMETRY IS MORE THAN TOPOLOGY

Classical differential geometry:

- an orientable manifold M , smooth functions, $C^\infty(M)$,
- differential algebra $\Omega(M)$, metric $g^{\mu\nu}$, Laplace operator Δ ,
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Problems are to calculate:

- the eigenvalues of the Dirac operator
- the invariants of the manifolds/structures

GEOMETRY AND THE HILBERT SPACES.

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Much of **classical** geometry can be encoded in terms of operators on a separable Hilbert space.

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- 5 dimension (growth of eigenvalues: $N(\Lambda) \sim \Lambda^d$),
- 6 integral (exotic traces) and other beasts...

THE GEOMETRY ACCORDING TO CONNES

THE SPECTRAL TRIPLE

Algebra \mathcal{A} , its faithful representation π on a Hilbert space \mathcal{H} , a selfadjoint unbounded operator D , satisfying several conditions:

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- 5 ...+ conditions of „analysis” type

THEOREM [CONNES]

If $\mathcal{A} = C^\infty(M)$, M a spin Riemannian compact manifold, $\mathcal{H} = L^2(S)$ (sections of spinor bundle) and D the Dirac operator on M then to $(\mathcal{A}, \mathcal{H}, D)$ is a spectral triple (with a real structure).

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A. Connes, *Noncommutative geometry and reality*, J. Math. Phys. 36, 6194, (1995)

COMMUTATIVE GEOMETRIES

which satisfy Connes' axioms are in 1:1 correspondence with Riemannian spin manifolds with a given spin structure and metric.

A. Connes, *On the spectral characterization of manifolds*, J. Noncom. Geom. 7, 1–82 (2013)

REMARK

Classical (real) spectral triples are *slightly* richer than spin geometries – as they describe (for example) geometries with torsion.

GENUINE NONCOMMUTATIVE REAL SPECTRAL TRIPLES

EXAMPLES OF REAL SPECTRAL GEOMETRIES

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HOW TO CONSTRUCT THEM?

There is **so far** no general method. Only examples.

REDISCOVERING THE SPIN STRUCTURE.

The question: What is the spin structure in NCG ?

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THEOREM (PASCHKE & SITARZ, LMP, 77, 3,(2006))

There are four inequivalent equivariant spin structures on the 2-dimensional noncommutative torus, with a unique choice of equivariant Dirac operator for each spin structure:

$$d_{\mu,\nu}^+ = \tau_\mu \mu + \tau_\nu \nu,$$

which satisfies the Hochschild cycle condition, provided that $\tau_\mu \tau_\nu^ \neq \tau_\mu^* \tau_\nu$. The spectrum of the equivariant Dirac operator depends on the spin structure.*

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More results followed (J-J. Venselaar).

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Problem: Why only flat or round geometries ?

ARE THERE ANY INTERESTING NC GEOMETRIES ?

A SOFTER VERSION OF *geometry*?

The facts:

- 1 for the examples of q -deformed algebras (Podleś spheres, $SU_q(2)$) - there are no spectral geometries **in the exact sense** – but – there are geometries in which some of the commutation relations are **satisfied up to compact operators**:

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Remark: Leads to nontrivial classical "triples".

RECENT EXAMPLES OF NEW NC GEOMETRIES

GEOMETRIES FROM NC CIRCLE BUNDLES

Take M a compact Riemannian spin manifold, on which S^1 acts freely and isometrically. Assume that the length of fibre is constant. **Aim:** express the Dirac operator on the total space using the Dirac on the base space and the $U(1)$ connection ω .
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CONFORMAL DEFORMATIONS OF NC TORI AND TORIC MANIFOLD

A family of conformally rescaled Dirac operators on the noncommutative 2-torus for which the Gauss-Bonnet formula holds:

$$D_h = hDh, \quad h^2 D^2 h^2,$$

where $h \in \mathcal{JC}^\infty(\mathbb{T}_\theta^2)J$, so it is in the commutant, $h > 0$, was introduced by Connes and Tretkoff, by M.Khalkhali et al, LD,AS. All good properties (Hochschild cocycle etc) hold.

RECENT EXAMPLES OF NEW NC GEOMETRIES

PARTIAL CONFORMAL DEFORMATIONS

If you take a torus with the metric $dx^2 + k^{-2}(x, y)dy^2$ (that is, for instance the usual „round“ torus embedded in \mathbb{R}^3) the Dirac operator is:

$$D = -i\sigma^1 \partial_x - i\sigma^2 \left(k \partial_y + \frac{1}{2} \partial_y(k) \right),$$

Same is possible with NC torus and the Gauss-Bonnet holds (LD+AS, Asymmetric noncommutative torus, SIGMA 11 (2015) 075-086).

These are examples of **new** spectral geometries that **do not** satisfy (or at least not in the obvious sense) the axioms of first-order condition.

NEW: REALITY TWISTED BY AN AUTOMORPHISM

Let A be a complex $*$ -algebra and let (H, π) be a (left) representation of A on a complex vector space H . A linear automorphism ν of H defines an algebra automorphism

$$\bar{\nu} : \text{End}(H) \rightarrow \text{End}(H), \quad \phi \mapsto \nu \circ \phi \circ \nu^{-1}.$$

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is an algebra map too, and hence it defines a new representation (H, π^ν) of A . The map ν is an isomorphism that intertwines (H, π) with (H, π^ν) .

We could also require that $\pi^\nu(a) \in \pi(A)$ so for faithful π the map $\bar{\nu}$ defines an (algebra) automorphism of A

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DEFINITION (TWISTED REAL SPECTRAL TRIPLE)

Let A be a $*$ -algebra, (H, π) a representation of A , D a linear operator on H , and let ν be a linear automorphism of H . We say that the triple (A, H, D) admits a ν -twisted real structure if there exists an anti-linear map $J : H \rightarrow H$ such that $J^2 = \epsilon \text{id}$, and, for all $a, b \in A$,

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$$\nu J\nu = J,$$

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If (A, H, D) admits a grading operator $\gamma : H \rightarrow H$:

$$\gamma^2 = \text{id}, \quad [\gamma, \pi(\mathbf{a})] = 0, \quad \gamma D = -D\gamma, \quad \nu^2 \gamma = \gamma \nu^2,$$

then the twisted real structure J is also required to satisfy

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In case of H being a Hilbert space the automorphism ν is also assumed to be densely defined and selfadjoint, with the requirement that $\bar{\nu}$ maps $\pi(A)$ into bounded operators.

The signs $\epsilon, \epsilon', \epsilon''$ determine the KO -dimension modulo 8 in the usual way and the operator J is antiunitary.

TWISTED REAL SPECTRAL TRIPLES

We shall say that a spectral triple admits a ν -twisted real structure, or simply that is a ν -twisted real spectral triple.

The commutant condition is called the *order-zero condition* and the one with the Dirac operator is called the *twisted order-one condition*. We shall call the modified condition the *the twisted ϵ' -condition*.

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REMARK

This is an **extension** not a **replacement**. In the case of $\nu = \text{id}$ we get the usual, well known, spectral triples.

THE FLUCTUATIONS OF THE DIRAC OPERATOR

Let Ω_D^1 be a bimodule of one forms:

$$\Omega_D^1 := \left\{ \sum_i \pi(a_i)[D, \pi(b_i)] \mid a_i, b_i \in A \right\}.$$

The standard fluctuation (= gauge transform) of a spectral triple (A, H, D) consist of

$$D \rightsquigarrow D + \alpha, \quad \alpha = \alpha^* \in \Omega_D^1.$$

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$$\Omega_D^1 := \left\{ \sum_i \pi(a_i)[D, \pi(b_i)] \mid a_i, b_i \in A \right\}.$$

The standard fluctuation (= gauge transform) of a spectral triple (A, H, D) consist of

$$D \rightsquigarrow D + \alpha, \quad \alpha = \alpha^* \in \Omega_D^1.$$

In case of a real spectral triple the fluctuated D is $D + \alpha + \epsilon' J\alpha J^{-1}$, where $\alpha + \epsilon' J\alpha J^{-1}$ is selfadjoint.

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For our case of ν -twisted real spectral triple we set the fluctuated Dirac operator D_α to be:

$$D_\alpha := D + \alpha + \epsilon' \nu J \alpha J^{-1} \nu,$$

with the requirement that $\alpha + \epsilon' \nu J \alpha J^{-1} \nu$ is selfadjoint.

FLUCTUATIONS

PROPOSITION

If (A, H, D) with $J \in \text{End}(H)$ is a ν -twisted real spectral triple, then (A, H, D_α) with (the same) J is also a ν -twisted real spectral triple.

If (A, H, D) is even with grading γ , then (A, H, D_α) is even with (the same) grading γ .

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The composition of twisted fluctuations is a twisted fluctuation.

PROOF

As a perturbation of D by a bounded selfadjoint operator, the fluctuated Dirac operator D_α is selfadjoint, has bounded commutators with $\pi(a) \in A$ and has compact resolvent.

We show that a fluctuation of the fluctuated Dirac operator is also a fluctuation. In other words, that

$$\Omega_{D_\alpha}^1 = \Omega_D^1, \quad \alpha \in \Omega_D^1.$$

PROOF (CONTINUED)

We compute:

$$\begin{aligned}[\nu J \alpha J^{-1} \nu, \pi(\mathbf{a})] &= \nu J \alpha J^{-1} \nu \pi(\mathbf{a}) - \pi(\mathbf{a}) \nu J \alpha J^{-1} \nu \\ &= \nu J \alpha J^{-1} \nu \pi(\mathbf{a}) - \nu \pi(\hat{\nu}^{-1}(\mathbf{a})) J \alpha J^{-1} \nu \\ &= \nu J \alpha J^{-1} \nu \pi(\mathbf{a}) - \nu J \alpha J^{-1} \pi(\hat{\nu}(\mathbf{a})) \nu \\ &= \nu J \alpha J^{-1} \nu \pi(\mathbf{a}) - \nu J \alpha J^{-1} \nu \pi(\mathbf{a}) \nu = 0.\end{aligned}$$

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We compute:

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Therefore for any $\alpha \in \Omega_D^1$ and $\mathbf{a} \in A$ we have:

$$[D_\alpha, \pi(\mathbf{a})] = [D, \pi(\mathbf{a})] + [\alpha, \pi(\mathbf{a})] \in \Omega_D^1.$$

PROOF (CONTINUED)

To finish the proof it remains only to check that D_α satisfies the twisted ϵ' -condition

$$D_\alpha J\nu = \epsilon'\nu JD_\alpha.$$

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Since D itself satisfies it, just check that

$$\begin{aligned}(\alpha + \epsilon' \nu J \alpha J^{-1} \nu) J \nu &= (\alpha J \nu + \epsilon' \nu J \alpha J^{-1} \nu J \nu) \\ &= \alpha J \nu + \epsilon' \nu J \alpha \\ &= \epsilon' \nu J (\alpha + \epsilon' J^{-1} \nu^{-1} \alpha J \nu) \\ &= \epsilon' \nu J (\alpha + \epsilon' \nu J \alpha J^{-1} \nu). \quad \square\end{aligned}$$

PROOF (CONTINUED)

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Since D itself satisfies it, just check that

$$\begin{aligned}(\alpha + \epsilon'\nu J\alpha J^{-1}\nu)J\nu &= (\alpha J\nu + \epsilon'\nu J\alpha J^{-1}\nu J\nu) \\ &= \alpha J\nu + \epsilon'\nu J\alpha \\ &= \epsilon'\nu J(\alpha + \epsilon'J^{-1}\nu^{-1}\alpha J\nu) \\ &= \epsilon'\nu J(\alpha + \epsilon'\nu J\alpha J^{-1}\nu). \quad \square\end{aligned}$$

Thus like in the usual case of the real spectral triples the twisted fluctuations form a semigroup.

EXAMPLE: CONFORMAL PERTURBATIONS

Let us assume that we have a real spectral triple (A, H, D, J) with reality operator J and fixed signs ϵ, ϵ' . Let $k \in \pi(A)$ be a positive and invertible bounded operator such that k^{-1} is also bounded, and let us denote by $k^J := Ad_J(k) = JkJ^{-1}$.

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PROPOSITION

If (A, H, D) with J is a real spectral triple, which satisfies order one condition, then for:

$$D_k = k^J D k^J, \quad \nu(h) = k^{-1} k^J h,$$

the triple (A, H, D_k) with J is a ν -twisted real spectral triple. If furthermore (A, H, D) is even with grading γ , then (A, H, D_k) is even with (the same) grading γ .

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PROOF

Since k and k^J are bounded operators it is clear that $\bar{\nu}$ sends bounded operators to bounded operators, and $\forall a \in A$:

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We show now that D_k satisfies the twisted order-one condition :

$$\begin{aligned} J\pi(b)J^{-1}[D_k, \pi(a)] &= J\pi(b)J^{-1}JkJ^{-1}[D, \pi(a)]JkJ^{-1} \\ &= k^J[D, \pi(a)]k^J J(k^{-2}\pi(b)k^2)J^{-1} = [D_k, \pi(a)]J\bar{\nu}^2(\pi(b))J^{-1}. \end{aligned}$$

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Next we check compatibility between J and ν :

$$\nu J\nu = k^{-1}JkJ^{-1}Jk^{-1}JkJ^{-1} = J.$$

EXAMPLE 1: CONFORMAL PERTURBATIONS

PROOF (CTD)

Finally, if $JD = \epsilon' DJ$ then for D_k we have:

$$JD_k = JK^J J^{-1} JDk^J = \epsilon' kDJk^J = \epsilon' k(k^J)^{-1} D_k(k^J)^{-1} kJ,$$

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REMARK

In the 'classical' case of a manifold M and (commutative) $A = C^\infty(M)$ with Ad_J being the complex conjugation, the conformal twists are always trivial as $JkJ^{-1} = k$ for a positive k and hence $\nu = \text{id}$.

THE (ν, ρ) TWISTING

DEFINITION $((\nu, \rho)$ -TWISTED ST)

We say that (A, H, D, J) is a (ν, ρ) -type twisted real spectral triple if:

- (1) for all $a \in A$, the commutators $[D, a]_\rho$ are bounded,
- (2) νJ preserves the domain of D ,
- (3) $DJ\nu = \epsilon' \nu J D$ and $\nu J \nu = J$ and $\nu^2 \gamma = \gamma \nu^2$,
- (4) the (ν, ρ) -twisted first-order condition holds:

$$[[D, a]_\rho, b]_{\rho \circ \nu^{-2}} = 0$$

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EXAMPLES

- (1) The (ν, id) -type spectral triple (untwisted) with twisted reality of [Brzezinski, Ciccola, Dabrowski, Sitarz]
- (2) The $(1, \rho)$ -type twisted real spectral triple of [Landi, Martinetti].

TWISTING AND UNTWISTING

THEOREM

Let A be a $*$ -algebra and $\tilde{\pi} : A \rightarrow B(H)$ be a $*$ -representation of A on a Hilbert space H . Let $J : H \rightarrow H$ be a \mathbb{C} -antilinear isometry such that $J^2 = \epsilon$ and that the zero order condition is satisfied. Let ρ be an algebra automorphism, and let ν be a bounded operator on H with the bounded inverse such that

- (a) ν implements an algebra automorphism $\hat{\nu}$ of A in representation $\tilde{\pi}$ and $\rho = \hat{\nu}^{-2}$, or
- (b) ν is a unitary operator such that ν^{-2} implements ρ in representation $\tilde{\pi}$

Let

$$\pi_\nu : A \rightarrow B(H), \quad a \mapsto \nu^{-1} \tilde{\pi}(a) \nu, \quad (1)$$

be the induced representation of A ...

TWISTING AND UNTWISTING

... and set

$$\pi = \begin{cases} \tilde{\pi}, & \text{in case (a),} \\ \pi_\nu, & \text{in case (b),} \end{cases}$$

so that π is always a $*$ -representation. Assume further that

$$\nu J \nu = J.$$

For an operator \tilde{D} on H , set

$$D = \nu \tilde{D} \nu,$$

Then:

- (1) (π, D, J, ν^2) satisfy conditions of a spectral triple with a ν^2 -twisted real structure if and only if $(\tilde{\pi}, \tilde{D}, J, \rho)$ satisfy conditions of real ρ -twisted spectral triple.

TWISTED AND UNTWISTED

We can summarise here three different kinds of twisted reality conditions obtained by the conformal twisting of a real spectral triple (A, H, π, D, J) in the following table:

$(A, H, \pi, k'Dk', J)$	$(A, H, \pi, kk'Dkk', J)$	(A, H, π, kDk, J)
spectral triple with the ν -twisted real structure and first-order condition	real ρ -twisted spectral triple	twisted spectral triple with real structure and untwisted first-order condition
$\nu = k^{-1}k'$	$\rho = \text{Ad}_{u^2}$	$\nu = kk'^{-1}$

Here $k = \pi(u) \in \pi(A)$, where $u \in A$ is invertible and such that k is positive with bounded inverse, $k' = JkJ^{-1}$ and we have $\nu JD = \epsilon'JD\nu$, and $\nu J\nu = J$ in the first and the third cases.

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THANK YOU !

REFERENCES

Twisted reality condition for Dirac operators

T.Brzeziński, N. Ciccoli, L.Dabrowski, AS

arXiv:1601.07404,

Math Phys Anal Geom (2016) 19:16

Twisted reality condition for spectral triple on two points

L.Dabrowski, AS

arXiv: 1605.03760,

Corfu Summer Institute 2015, PoS(CORFU2015)093

On twisted reality conditions

T.Brzeziński, L.Dabrowski, AS

arXiv: 1804.07005,

Lett Math Phys (2018) 108:1323–1340